

Partial Differential Equation Solving, vol 2.

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Monte Carlo methods,
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There will be no lectures and class on 19th of May

Dirichlet conditions: expected number of steps

⇒ find the function $u(x_1, x_2, \dots, x_k)$ such that it fulfils the Laplace equation:

$$\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_k^2} = 0, \quad (x_1, x_2, \dots, x_k) \in D \subset \mathbb{R}^k$$

In the domain D , on the the $\Gamma(D)$ the u function is given by:

$$U(x_1, x_2, \dots, x_k) = f(x_1, x_2, \dots, x_k), \quad (x_1, x_2, \dots, x_k) \in \Gamma(D)$$

⇒ Now lets assume that the domain D is a hyperball:

$$0 \leq \sum_{i=1}^k x_i^2 \leq r^2, \quad r = \text{const}$$

⇒ Now $\pi_\nu(x_1, x_2, \dots, x_k)$ is a probability that a particle starting from (x_1, x_2, \dots, x_k) will end up on the edge after ν steps. The $\kappa(x_1, x_2, \dots, x_k)$ is the estimated number of steps for this trajectory.

$$\pi_0 = \begin{cases} 1, & (x_1, x_2, \dots, x_k) \in \Gamma(D) \\ 0, & (x_1, x_2, \dots, x_k) \in D \end{cases} \quad (1)$$

$$\pi_\nu = \frac{1}{2k} \sum_{l=1}^{\nu} \pi_{\nu-l}(x_1^l, x_2^l, \dots, x_k^l)$$

Dirichlet conditions: expected number of steps

⇒ From Eq. 1 and 2 one gets:

$$\kappa(x_1, x_2, \dots, x_k) = \sum_{\nu=1}^{\infty} \nu \pi_{\nu}(x_1, x_2, \dots, x_k)$$

one gets:

$$\begin{aligned} \kappa(x_1, x_2, \dots, x_k) &= \frac{1}{2k} \sum_{\nu=1}^{\infty} \left[\nu \sum_{l=1}^k \pi_{\nu-1}(x_1, x_2, \dots, x_k) \right] \\ &= \frac{1}{2k} \sum_{\nu=1}^{\infty} \left[(\nu-1) \sum_{l=1}^k \pi_{\nu-1}(x_{1'}, x_{2'}, \dots, x_{k'}) \right] + \frac{1}{2k} \sum_{\nu=1}^{\infty} \sum_{l=1}^k \pi_{\nu-1}(x_{1'}, x_{2'}, \dots, x_{k'}) \end{aligned}$$

⇒ From which we get:

$$\kappa(x_1, x_2, \dots, x_k) = \frac{1}{2k} \sum_{l=1}^k \kappa(x_{1'}, x_{2'}, \dots, x_{k'}) + 1$$

⇒ Now this is equivalent of the Poisson differential equation:

$$\frac{\partial^2 \kappa}{\partial x_1^2} + \frac{\partial^2 \kappa}{\partial x_2^2} + \dots + \frac{\partial^2 \kappa}{\partial x_k^2} = -2k, \text{ b. con. } \kappa(x_1, x_2, \dots, x_k) = 0, (x_1, x_2, \dots, x_k) \in \Gamma(D)$$

Dirichlet conditions: expected number of steps

\Rightarrow From previous equation: $\kappa(x_1, x_2, \dots, x_k) = \psi(x_1, x_2, \dots, x_k) - \sum_{i=1}^k x_i^2$ we get the for the ψ function the Laplace equation:

$$\frac{\partial^2 \psi}{\partial x_1^2} + \frac{\partial^2 \psi}{\partial x_2^2} + \dots + \frac{\partial^2 \psi}{\partial x_k^2} = 0$$

because on the border ($\Gamma(D)$):

$$\psi(x_1, x_2, \dots, x_k) = r^2 = \text{const}$$

so also inside the D : $\psi(x_1, x_2, \dots, x_k) = r^2 = \text{const} \Rightarrow$ From which we can estimate the number steps in the random walk:

$$\kappa(x_1, x_2, \dots, x_k) = r^2 - \sum_{i=1}^k x_i^2 \leq r^2$$

Important conclusion:

The expected number of steps in the random walk (the time of walk) from the point (x_1, x_2, \dots, x_k) till the edge of the domain can be estimated by r number (the LINEAR! size). It is completely independent of the k !

Dirichlet conditions as linear system

⇒ In the discrete form we can write the Dirichlet conditions as (2-dim case):

$$u(x, y) = \frac{1}{4} [u(x-1, y) + u(x+1, y) + u(x, y-1) + u(x, y+1)], \quad (x, y) \in D$$

$$u(x, y) = f(x, y), \quad (x, y) \in \Gamma(D)$$

⇒ Now we can order the grid $((x, y) \in D \cup \Gamma(D))$, we can represent the above equations as a linear system:

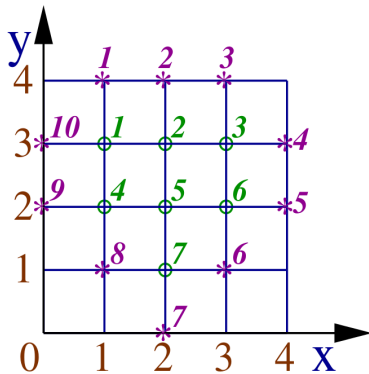
$$u_i = a_i + \sum_{j=1}^n h_{ij} u_j, \quad i = 1, 2, \dots, n$$

The trick:

So to solve a differential equation with Dirichlet boundary condition we can use all the methods of solving linear equation systems such as Neumann-Ulam or Wassow.

Dirichlet conditions as linear system - example

- To do this we act as following: we number separately the points inside the D domain and on the border $\Gamma(D)$.
- We write for each point inside the domain the Laplace equation as system of linear equations:



$$\begin{array}{rcll}
 u_1 & -u_2/4 & -u_4/4 & = (f_1 + f_{10})/4 \\
 -u_1/4 & u_2 - u_3/4 & -u_5/4 & = (f_2)/4 \\
 & -u_2/4 & u_3 & \\
 -u_1/4 & & u_4 - u_5/4 & -u_6/4 = (f_3 + f_4)/4 \\
 -u_1/4 & & -u_4/4 & u_5 & -u_6/4 - u_7/4 & = (f_8 + f_9)/4 \\
 & & & & & = 0 \\
 & -u_3/4 & & -u_5/4 & u_6 & = (f_5 + f_6)/4 \\
 & & & -u_5/4 & & u_7 & = (f_5 + f_6)/4
 \end{array}$$

Dirichlet conditions as linear system - example

⇒ The above equation we can transform the above equation into the iterative representation:

$$\vec{u} = \vec{a} + \mathbf{H}\vec{u}$$

where $\vec{u} = (u_1, u_2, \dots, u_7)$ is the vector which represent the values of the function inside the D domain, \vec{a} is the linear combinations of the boundary values. In our example:

$$\mathbf{H} = \begin{pmatrix} 0 & \frac{1}{4} & 0 & \frac{1}{4} & 0 & 0 & 0 \\ \frac{1}{4} & 0 & \frac{1}{4} & 0 & \frac{1}{4} & 0 & 0 \\ 0 & \frac{1}{4} & 0 & 0 & 0 & \frac{1}{4} & 0 \\ \frac{1}{4} & 0 & 0 & 0 & \frac{1}{4} & 0 & 0 \\ 0 & \frac{1}{4} & 0 & \frac{1}{4} & 0 & \frac{1}{4} & \frac{1}{4} \\ 0 & 0 & \frac{1}{4} & 0 & \frac{1}{4} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{4} & 0 & 0 \end{pmatrix}$$

⇒ To find the solution to aka \vec{u} one can use the methods we already know: Neumann-Ulam and Wasow, etc.

⇒ There are tricks and tips one can use to make this problem faster as each of the entry is $\frac{1}{4}$.

Neumann-Ulam method

- ⇒ We put the particle in (x, y) .
- ⇒ We observe the trajectory of the particle until it reaches the boundary. Point P_k is the last point before hitting the boundary.
- ⇒ For each trajectory we assign a value that is the arithmetical mean of the boundary points that are neighbours of the point P_k .
- ⇒ We repeat the above n times and calculate the mean.
- ⇒ The example solution for 20 trajectories:

$$u(2, 2) = 1.0500 \pm 0.2756$$

- ⇒ E 10.1 Solve the above linear system using the Neumann-Ulam method for an assumed boundary conditions.

Dual Wasow method

- ⇒ We choose the boundary conditions with arbitrary chosen probability p.d.f. $p(Q)$ the starting point.
- ⇒ We choose with equal probability the point inside D where the particle goes.
- ⇒ With equal probability we choose the next positions and so on until the particle hits the boundary in the point Q' .
- ⇒ We count all trajectories $N((x_1, x_2, x_3, \dots, x_k))$ that that have passed the point $(x_1, x_2, x_3, \dots, x_k)$.
- ⇒ For the point (x_1, x_2, \dots, x_k) we calculate:

$$w(x_1, x_2, \dots, x_k) = \frac{1}{2k} N(x_1, x_2, \dots, x_k) \frac{f(Q)}{p(Q)}$$

- ⇒ The above steps we repeat N times.
- ⇒ After that we take the arithmetic mean of w .

Random walk with different step size

⇒ If $u(x, y)$ is a harmonic function that obeys the Laplace equation and $S_r(x, y)$ is a circle in with the middle point (x, y) and radius r . Then a theorem states:

$$S_r(x, y) = \frac{1}{2\pi} \int_0^{2\pi} u(x + r \cos \phi, y + r \sin \phi) d\phi$$

⇒ The above is true for in all the dimensions.

⇒ The E.Muller method:

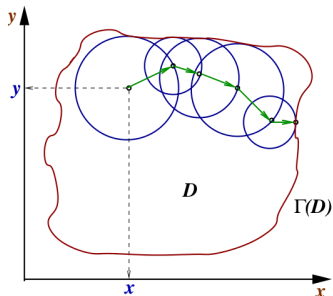
- At the begging we set the point in the initial point: (x_1, x_2, \dots, x_k) .
- We construct a k dimensional sphere with center (x_1, x_2, \dots, x_k) and radius r . The r has to be chosen in a way that the whole is inside the D : $S_r(\vec{x}) \in D$. We choose a random point from $\mathcal{U}(0, 2\pi)$ on the sphere which is our new point.
- We stop the walk when the point is on $\Gamma(D)$.

⇒ We repeat this N times.

⇒ The final result if the arithmetical mean of all trajectories and is equal of the $u(x_1, x_2, \dots, x_k)$.

Muller method

- ⇒ The method is faster the faster the particle reaches the edge.
- ⇒ In order to do so we choose the radius that it is the maximal one that allows the sphere to be inside the domain D .



- ⇒ There is a problem!!!! The probability that we choose a point on the edge is 0!!!!
- ⇒ An approximation has to be made: we choose a small number δ and we consider that the particle reached the border when the distance is with δ .
- ⇒ We can always choose the δ such that the estimator error of function is smaller then a given ϵ .

Muller method, example

⇒ An example solution in the

Other boundary conditions

⇒ Find the solution to the Laplace equation:

$$\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_k^2} = 0, \quad (x_1, x_2, \dots, x_k) \in D \subset \mathbb{R}^k$$

inside the D domain if on the edge $\Gamma(D)$ the function fulfils the equation:

$$f(x_1, x_2, \dots, x_k) \frac{\partial u(x_1, x_2, \dots, x_k)}{\partial n} + g(x_1, x_2, \dots, x_k) u(x_1, x_2, \dots, x_k) = h(x_1, x_2, \dots, x_k)$$

where $\frac{\partial u(x_1, x_2, \dots, x_k)}{\partial n}$ is there derivative in the direction of normal to the $\Gamma(D)$ in the direction inside D .

⇒ The cases:

- $f = 0$. ⇒ Dirichlet boundary condition (1st class condition).
- $g = 0$. ⇒ Neumann boundary condition (2nd class condition).
- others. ⇒ General case (3rd class condition).

Other boundary conditions

⇒ In 2-dim:

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, \quad (x, y) \in D \subset \mathbb{R}^2$$

with the boundary condition:

$$f(x, y) \frac{\partial u(x, y)}{\partial n} + g(x, y) u(x, y) = h(x, y), \quad (x, y) \in \Gamma(D)$$

⇒ And the discrete differential equation:

$$u(x, y) = \frac{1}{4} [u(x - h, y) + u(x + h, y) + u(x, y - h) + u(x, y + h)]$$

Reminder:

If at moment t the point is in (x, y) then in the $t + 1$ time the particle moves with equal probability to one of the following points: $(x - h, y)$, $(x + h, y)$, $(x, y - h)$, $(x, y + h)$.

Random walk for boundary points

⇒ The boundary point Q has only one internal neighbour point P .

- If the normal is parallel to the grid axis in the point Q :

$$f(Q) \frac{u(P) - u(Q)}{h} + g(Q)u(Q) = h(Q)$$

- Solving the above to get $u(Q)$ we get:

$$u(Q) = \frac{f(Q)u(P)}{f(Q) - hg(Q)} - \frac{h(Q)}{f(Q) - hg(Q)}$$

- To help we assign a temporary values:

$$\phi(Q) = \frac{f(Q)}{p[f(Q) - hg(Q)]}, \quad \psi(Q) = -h \frac{h(Q)}{(1-p)[f(Q) - hg(Q)]}$$

⇒ So:

$$u(Q) = p\phi(Q)u(P) + (1-p)\psi(Q)$$

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⇒ So:

$$u(Q) = p\phi(Q)u(P) + (1-p)\psi(Q)$$

⇒ Interpretation: $u(Q)$ can be seen that with probability p it is equal $\phi(Q)u(P)$ and with provability $(1-p)$ is equal to $\psi(Q)$.

Backup