

# Markov Chain MC

Marcin Chrzęszcz  
mchrzasz@cern.ch



Experimental Methods in Particle Physics,  
11 October, 2017

## Trivial example

⇒ Lets start with a TRIVIAL example: we want to calculate  $S = A + B$ .

We can rewrite it in:

$$S = p \frac{A}{p} + (1 - p) \frac{B}{1 - p}$$

and one can interpret the sum as expected value of:

$$W = \begin{cases} \frac{A}{p} & \text{with propability } p \\ \frac{A}{1-p} & \text{with propability } 1 - p \end{cases}$$

⇒ The algorithm:

- We generate a random variable  $W$  and calculate:

$$\hat{S} = \frac{1}{N} \sum_{i=1}^N W_i$$

⇒ The  $\hat{S}$  is an unbiased estimator of  $S$ .

## Trivial example2

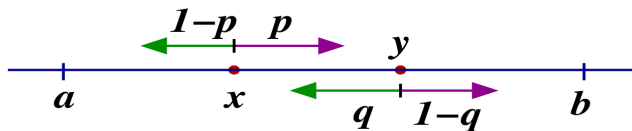
- Lets say we have a linear equation system:

$$\begin{aligned}X &= pY + (1 - p)A \\ Y &= qX + (1 - q)B\end{aligned}$$

- We know  $A, B, p, q$ ;  $X$  and  $Y$  are meant to be determined.
- Algorithm:
  1. We choose first element of the first equation with probability  $p$  and second with probability  $1 - p$ .
  2. If we choose the second one, the outcome of this MCMC is  $W = A$ .
  3. If we choose the first we go to second equation and choose the first element with probability  $q$  and the second with  $1 - q$ .
  4. If we choose the second one, the outcome of this MCMC is  $W = B$ .
  5. If we choose the first we go to the first equation back again.
  6. We repeat the procedure.
- We can estimate the solution of this system:

$$\hat{X} = \frac{1}{N} \sum_{i=1}^N W_i \quad \hat{\sigma}_X = \frac{1}{\sqrt{N-1}} \sqrt{\frac{1}{N} \sum_{i=1}^N W_i^2 - \hat{X}^2}$$

# Random walk



- ⇒ We are in the point  $x$  and we walk accordingly to the following rules:
- From point  $x$  we walk with probability  $p$  to point  $y$  or with  $1 - p$  to  $a$ .
  - From point  $y$  we walk with probability  $q$  to point  $x$  and with  $1 - Q$  to  $b$ .
  - The walks ends when you end up in  $a$  or  $b$ .
  - You get a "reward"  $A$  if you end up in point  $a$  and  $B$  if you end up in  $b$ .
  - $X$  is expected "reward" when you start the walk from  $x$ ,  $Y$  when you start from  $y$ .
- ⇒ The algorithm above is so-called random walk on the set  $\{a, x, y, b\}$
- ⇒ The described walked can solve the linear equation system that we discussed above.

# Markov Chain MC

- Consider a finite (or Countable set) possible states:  $S_1, S_2, \dots$
- The  $X_t$  is the state of the system in the time  $t$
- We are looking at discrete time steps: 1, 2, 3, ...
- The conditional probability is defined as:

$$P(X_t = S_j | X_{t-1} = S_{j-1}, \dots, X_1 = S_1)$$

- The Markov chain is then if the probability depends only on previous step.

$$P(X_t = S_j | X_{t-1} = S_{j-1}, \dots, X_1 = S_1) = P(X_t = S_j | X_{t-1} = S_{j-1})$$

- For this reason MCMC is also known as drunk sailor walk.
- Very powerful method. Used to solve linear eq. systems, invert matrix, solve differential equations, etc.
- Also used in physics problems: Brown motions, diffusion, etc.

# Linear equations system

⇒ Lets start from a linear equation system:

$$\mathbf{A}\vec{x} = \vec{b}, \quad \det \mathbf{A} \neq 0,$$

where  $\mathbf{A} = (a_{ij}, i, j = 1, 2, \dots, n)$  -matrix,  $\vec{b} = (b_1, b_2, \dots, b_n)$ -vector,  
 $\vec{x} = (x_1, x_2, \dots, x_n)$  - vector of unknowns.

⇒ The solution we mark as  $\vec{x}^0 = (x_1^0, x_2^0, \dots, x_n^0)$

⇒ The above system can be transformed into the iterative representation:

$$\vec{x} = \vec{a} + \mathbf{H}\vec{x}$$

where  $\mathbf{H}$  is a matrix,  $\vec{a}$  is a vector.

⇒ We assume that the matrix norm:

$$\|\mathbf{H}\| = \max_{1 \leq i \leq n} \sum_{j=1}^n |h_{ij}| < 1$$

# Linear equations system

⇒ Lets start from a linear equation system:

$$\mathbf{A}\vec{x} = \vec{b}, \quad \det \mathbf{A} \neq 0,$$

where  $\mathbf{A} = (a_{ij}, i, j = 1, 2, \dots, n)$  -matrix,  $\vec{b} = (b_1, b_2, \dots, b_n)$ -vector,  
 $\vec{x} = (x_1, x_2, \dots, x_n)$  - vector of unknowns.

⇒ The solution we mark as  $\vec{x}^0 = (x_1^0, x_2^0, \dots, x_n^0)$

⇒ The above system can be transformed into the iterative representation:

$$\vec{x} = \vec{a} + \mathbf{H}\vec{x}$$

where  $\mathbf{H}$  is a matrix,  $\vec{a}$  is a vector.

⇒ We assume that the matrix norm:

$$\|\mathbf{H}\| = \max_{1 \leq i \leq n} \sum_{j=1}^n |h_{ij}| < 1$$

⇒ We can always change transform every system to the iteration form:  $\mathbf{A} = \mathbf{V} - \mathbf{W}$ .

$$(\mathbf{V} - \mathbf{W})\vec{x} = \vec{b} \quad \mapsto \quad \vec{x} = \mathbf{V}^{-1}\vec{b} + \mathbf{V}^{-1}\mathbf{W}\vec{x}$$

# Linear equations system

⇒ Now we further modify the equation system:

$$\vec{x} = \vec{a} + \mathbf{H}\vec{x} \Rightarrow (\mathbf{I} - \mathbf{H})\vec{x} = \vec{a}$$

where  $\mathbf{I} = \delta_{ij}$  - unit matrix,  $\delta_{ij}$  is the Kronecker delta.

⇒ What one can do is to represent the solution in terms of Neumann series:

$$\vec{x}^0 = (\mathbf{I} - \mathbf{H})^{-1}\vec{a} = \vec{a} + \mathbf{H}\vec{a} + \mathbf{H}^2\vec{a} + \mathbf{H}^3\vec{a} + \dots$$

⇒ So for the  $i^{th}$  component we have:

$$\begin{aligned} x_i^0 &= a_i + \sum_{j=1}^n h_{ij}a_j + \sum_{j_1=1}^n \sum_{j_2=1}^n h_{ij_1}h_{j_1j_2}a_{j_2} \\ &+ \dots + \sum_{j_1=1}^n \dots \sum_{j_n=1}^n h_{ij_1} \dots h_{j_{n-1}j_n}a_{j_n} \end{aligned}$$

⇒ We will construct a probabilistic interpretation using MCMC and then we show that the expected value is equal to the above formula.



# Neumann-Ulam method

- To do so we add to our matrix an additional column of the matrix:

$$h_{i,0} = 1 - \sum_{j=1}^n h_{ij} > 0$$

- The system has states:  $\{0, 1, 2, \dots, n\}$
- State at  $t$  time is denoted as  $i_t$  ( $i_t = 0, 1, 2, \dots, n; t = 0, 1, \dots$ )
- We make a random walk accordingly to the following rules:
  - At the beginning of the walk ( $t = 0$ ) we are at  $i_0$ .
  - In the  $t$  moment we are in the  $i_t$  position then in  $t + 1$  time stamp we move to state  $i_{t+1}$  with the probability  $h_{i_t i_{t+1}}$ .
  - We stop walking if we are in state 0.
- The path  $\gamma = (i_0, i_1, i_2, \dots, i_k, 0)$  is called trajectory.
- For each trajectory we assign a number:

$$X(\gamma) = X(i_0, i_1, i_2, \dots, i_k, 0) = \frac{a_{i_k}}{h_{i_k 0}}$$

# Neumann-Ulam method

- ⇒ The  $X(\gamma)$  variable is a random variable from:  $\{a_1/h_{1,0}, a_2/h_{2,0}, \dots, a_n/h_{n,0}\}$ .  
The probability that  $X(\gamma) = a_j/h_{j,0}$  is equal to the probability that the last non zero state of the  $\gamma$  trajectory is  $j$ .
- ⇒ The expected value of the  $X(\gamma)$  trajectory if the trajectory begins from  $i_0 = s$  is:

$$E\{X(\gamma)|i_0 = s\} = \sum_{k=0}^{\infty} \sum_{\{\gamma_k\}} X(\gamma)P(\gamma)$$

where  $\gamma_k$  is a trajectory of length  $k$ , which starts in  $i_0 = s$  and  $P(\gamma)$  is the probability of occurrence of this trajectory. ⇒ Yes you guess it lets do Taylor expansion:

$$E\{X(\gamma)|i_0 = s\} = \sum_{\gamma_0} X(\gamma)P(\gamma) + \sum_{\gamma_1} X(\gamma)P(\gamma) + \dots + \sum_{\gamma_k} X(\gamma)P(\gamma)$$

- ⇒ Now let's examine the elements of the above series.

# Neumann-Ulam method

$\{\gamma_0\}$ : One trajectory:  $\gamma_0 = (i_0 = s|0)$ ,  $P(\gamma_0) = h_{s,0}$  and  $X(\gamma_0) = a_s/h_{s,0}$ . So:

$$\sum_{\gamma_0} X(\gamma)P(\gamma) = \frac{a_s}{h_{s,0}} h_{s,0} = a_s$$

$\{\gamma_1\}$ : Trajectories:  $\gamma_1 = (i_0 = s, i_1|0)$ ,  $i_1 \neq 0$ ,  $P(\gamma_1) = P(s, i_1, 0) = h_{s,i_1} h_{i_1,0}$  and  $X(\gamma_1) = a_{i_1}/h_{i_1,0}$ . So:

$$\sum_{\gamma_1} X(\gamma)P(\gamma) = \sum_{i_1=1}^n \frac{a_{i_1}}{h_{i_1,0}} h_{s,i_1} h_{i_1,0} = \sum_{i_1=1}^n h_{s,i_1} a_{i_1}$$

$\{\gamma_2\}$ : Trajectories:  $\gamma_2 = (i_0 = s, i_1, i_2|0)$ ,  $i_1, i_2 \neq 0$ ,  
 $P(\gamma_2) = P(s, i_1, i_2, 0) = h_{s,i_1} h_{i_1,i_2} h_{i_2,0}$  and  $X(\gamma_2) = a_{i_2}/h_{i_2,0}$ . So:

$$\sum_{\gamma_2} X(\gamma)P(\gamma) = \sum_{i_1=1}^n \sum_{i_2=1}^n \frac{a_{i_2}}{h_{i_2,0}} h_{s,i_1} h_{i_1,i_2} h_{i_2,0} = \sum_{i_1=1}^n \sum_{i_2=1}^n h_{s,i_1} h_{i_1,i_2} a_{i_2}$$

⇒ etc...

# Neumann-Ulam method

⇒ After summing up:

$$\begin{aligned} E\{X(\gamma)|i_0 = s\} &= a_s + \sum_{i_1=1}^n h_{s,i_1} a_{i_1} + \sum_{i_1=1}^n \sum_{i_2=1}^n h_{s,i_1} h_{i_1,i_2} a_{i_2} + \dots \\ &+ \sum_{i_1=1}^n \sum_{i_2=1}^n \dots \sum_{i_k=1}^n h_{s,i_1} h_{i_1,i_2} \dots h_{i_{k-1},i_k} a_{i_k} + \dots \end{aligned}$$

⇒ If you compare this expression with the Neumann series we will they are the same so:

$$x_i^0 = E\{X(\gamma)|i_0 = i\}$$

To sum up:

We have proven that solving a linear system can be represented by an expectation value of the random variable  $X(\gamma)$ . The error is computed using standard deviation equation.

# Neumann-Ulam method

- For example lets try to solve this equation system:

$$\vec{x} = \begin{pmatrix} 1.5 \\ -1.0 \\ 0.7 \end{pmatrix} + \begin{pmatrix} 0.2 & 0.3 & 0.1 \\ 0.4 & 0.3 & 0.2 \\ 0.3 & 0.1 & 0.1 \end{pmatrix} \vec{x}$$

- The solution is  $\vec{x}_0 = (2.154303, 0.237389, 1.522255)$ .
- The propability matrix  $h_{ij}$  has the shape:

$i/j$	1	2	3	0
1	0.2	0.3	0.1	0.4
2	0.4	0.3	0.2	0.1
3	0.3	0.1	0.1	0.5

- An example solution:

```
mchraszcz-ThinkPad-W530% ./mark.x 1 1000000  
2.15625
```

# Neumann-Ulam dual method

- The problem with Neumann-Ulam method is that you need to repeat it for each of the coordinates of the  $\vec{x}_0$  vector.
- The dual method calculates the whole  $\vec{x}_0$  vector.
- The algorithm:
  - On the indexes:  $\{0, 1, \dots, n\}$  we set a probability distribution:  $q_1, q_2, \dots, q_n, q_i > 0$  and  $\sum_{i=1}^n q_i = 1$ .
  - The starting point we select from  $q_i$  distribution.
  - If in  $t$  time we are in  $i_t$  state then with probability  $p(i_{t+1}|i_t) = h_{i_{t+1}, i_t}$  in  $t + 1$  we will be in state  $i_{t+1}$ . For  $i_{t+1} = 0$  we define the probability:  $h_{0, i_t} = 1 - \sum_{j=1}^n h_{j, i_t}$ . Here we also assume that  $h_{j, i_t} > 0$ .
  - NOTE: there the matrix is transposed compared to previous method:  $H^T$ .
  - Again we end our walk when we are at state 0.
  - For the trajectory:  $\gamma = (i_0, i_1, \dots, i_k, 0)$ , we assign the vector:

$$\vec{Y}(\gamma) = \frac{a_{i_0}}{q_{i_0} p(0|i_k)} \hat{e}_{i_k} \in \mathcal{R}^n$$

- The solution will be :  $\vec{x}^0 = \frac{1}{N} \sum \vec{Y}(\gamma)$

# Neumann-Ulam dual method

- Let's try to solve the equation system:

$$\vec{x} = \begin{pmatrix} 1.5 \\ -1.0 \\ 0.7 \end{pmatrix} + \begin{pmatrix} 0.2 & 0.3 & 0.1 \\ 0.4 & 0.3 & 0.2 \\ 0.1 & 0.1 & 0.1 \end{pmatrix} \vec{x}$$

- The solution is:  $\vec{x}_0 = (2.0, 0.0, 1.0)$ .
- Let's put the initial probability as constant:

$$q_1 = q_2 = q_3 = \frac{1}{3}$$

- The propability matrix  $h_{ij}$  has the shape:

$i/j$	1	2	3	4
1	0.2	0.4	0.1	0.3
2	0.3	0.3	0.1	0.3
3	0.1	0.2	0.1	0.6

- An example solution:

```
mchrzasz-ThinkPad-W530% ./mark2.x 1000000  
1.9943 0.001806 1.00267
```

# Generalization, the algorithm

- ⇒ We set the  $P$  matrix in a arbitrary way.
- ⇒ If in the  $t$  moment the point is in the  $i_t$  state, then with the probability  $p_{i_t, i_{t+1}}$  he can go to  $i_{t+1}$  state.
- ⇒ We stop the walk once we reach 0.
- ⇒ For the given trajectory we assign the value:  $X(\gamma_k)$
- ⇒ We repeat the procedure  $N$  times and take the mean and RMS.
- ⇒ We repeat this also for every of the  $\vec{x}^0$  vector components.



# Partial differential equations, intro

⇒ Let's say we are want to describe a point that walks on the  $\mathbb{R}$  axis:

- At the beginning ( $t = 0$ ) the particle is at  $x = 0$
- If in the  $t$  the particle is in the  $x$  then in the time  $t + 1$  it walks to  $x + 1$  with the known probability  $p$  and to the point  $x - 1$  with the probability  $q = 1 - p$ .
- The moves are independent.

⇒ So let's try to described the motion of the particle.

⇒ The solution is clearly a probabilistic problem. Let  $\nu(x, t)$  be a probability that at time  $t$  particle is in position  $x$ . We get the following equation:

$$\nu(x, t + 1) = p\nu(x - 1, t) + q\nu(x + 1, t)$$

with the initial conditions:

$$\nu(0, 0) = 1, \quad \nu(x, 0) = 0 \text{ if } x \neq 0.$$

⇒ The above functions describes the whole system (every  $(t, x)$  point).

# Partial differential equations, intro

⇒ Now in differential equation language we would say that the particle walks in steps of  $\Delta x$  in times:  $k\Delta t$ ,  $k = 1, 2, 3, \dots$ :

$$\nu(x, t + \Delta t) = p\nu(x - \Delta x, t) + q\nu(x + \Delta x, t).$$

⇒ To solve this equation we need to expand the  $\nu(x, t)$  function in the Taylor series:

$$\begin{aligned}\nu(x, t) + \frac{\partial \nu(x, t)}{\partial t} \Delta t &= p\nu(x, t) - p \frac{\partial \nu(x, t)}{\partial x} \Delta x + \frac{1}{2} p \frac{\partial^2 \nu(x, t)}{\partial x^2} (\Delta x)^2 \\ &+ q\nu(x, t) + q \frac{\partial \nu(x, t)}{\partial x} \Delta x + \frac{1}{2} q \frac{\partial^2 \nu(x, t)}{\partial x^2} (\Delta x)^2\end{aligned}$$

⇒ From which we get:

$$\frac{\partial \nu(x, t)}{\partial t} \Delta t = -(p - q) \frac{\partial \nu(x, t)}{\partial x} \Delta x + \frac{1}{2} \frac{\partial^2 \nu(x, t)}{\partial x^2} (\Delta x)^2$$

⇒ Now We divide the equation by  $\Delta t$  and take the  $\Delta t \rightarrow 0$ :

$$(p - q) \frac{\Delta x}{\Delta t} \rightarrow 2c, \quad \frac{(\Delta x)^2}{\Delta t} \rightarrow 2D,$$

⇒ We get the Fokker-Planck equation for the diffusion with current:

$$\frac{\partial \nu(x, t)}{\partial t} = -2c \frac{\partial \nu(x, t)}{\partial x} + D \frac{\partial^2 \nu(x, t)}{\partial x^2}$$

⇒ The  $D$  is the diffusion coefficient,  $c$  is the speed of current. For  $c = 0$  it is a symmetric distribution.

# Laplace equation, Dirichlet boundary conditions

- ⇒ The aforementioned example shows the way to solve the partial differential equation using Markov Chain MC.
- ⇒ We will see how different classes of partial differential equations can be approximated with a Markov Chain MC, whose expectation value is the solution of the equation.
- ⇒ The Laplace equation:

$$\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} + \dots + \frac{\partial^2 u}{\partial x_k^2} = 0$$

The  $u(x_1, x_2, \dots, x_k)$  function that is a solution of above equation we call harmonic function. If one knows the values of the harmonic function on the edges  $\Gamma(D)$  of the  $D$  domain one can solve the equation.

## The Dirichlet boundary conditions:

Find the values of  $u(x_1, x_2, \dots, x_k)$  inside the  $D$  domain knowing the values of the edge are given with a function:

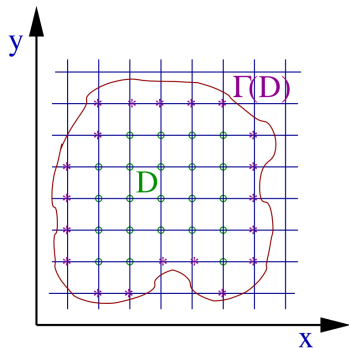
$$u(x_1, x_2, \dots, x_k) = f(x_1, x_2, \dots, x_k) \in \Gamma(D)$$

- ⇒ Now I am lazy so I put  $k = 2$  but it's the same for all  $k$ !

# Laplace equation, Dirichlet boundary conditions

⇒ We will put the Dirichlet boundary condition as a discrete condition:

- The domain  $D$  we put a lattice with distance  $h$ .
- Some points we treat as inside (denoted with circles). Their form a set denoted  $D^*$ .
- The other points we consider as the boundary points and they form a set  $\Gamma(D)$ .



⇒ We express the second derivatives with the discrete form:

$$\frac{\frac{u(x+h)-u(x)}{h} - \frac{u(x)-u(x-h)}{h}}{h} = \frac{u(x+h) - 2u(x) + u(x-h)}{h^2}$$

⇒ Now we choose the units so  $h = 1$ .

# Laplace equation, Dirichlet boundary conditions

The Dirichlet condition in the discrete form:

Find the  $u^*$  function which obeys the differential equation:

$$U^*(x, y) = \frac{1}{4} [u^*(x-1, y) + u^*(x+1, y) + u^*(x, y-1) + u^*(x, y+1)]$$

in all points  $(x, y) \in D^*$  with the condition:

$$u^*(x, y) = f^*(x, y), \quad (x, y) \in \Gamma(D^*)$$

where  $f^*(x, y)$  is the discrete equivalent of  $f(x, y)$  function.

⇒ We consider a random walk over the lattice  $D^* \cup \Gamma(D^*)$ .

- In the  $t = 0$  we are in some point  $(\xi, \eta) \in D^*$
- If at the  $t$  the particle is in  $(x, y)$  then at  $t + 1$  it can go with equal probability to any of the four neighbour lattices:  $(x-1, y)$ ,  $(x+1, y)$ ,  $(x, y-1)$ ,  $(x, y+1)$ .
- If the particle at some moment gets to the edge  $\Gamma(D^*)$  then the walk is terminated.
- For the particle trajectory we assign the value of:  $\nu(\xi, \eta) = f^*(x, y)$ , where  $(x, y) \in \Gamma(D^*)$ .

# Laplace equation, Dirichlet boundary conditions

⇒ Let  $p_{\xi,\eta}(x, y)$  be the probability of particle walk that starting in  $(\xi, \eta)$  to end the walk in  $(x, y)$ .

⇒ The possibilities:

1. The point  $(\xi, \eta) \in \Gamma(D^*)$ . Then:

$$p_{\xi,\eta}(x, y) = \begin{cases} 1, & (x, y) = \xi, \eta \\ 0, & (x, y) \neq \xi, \eta \end{cases} \quad (1)$$

2. The point  $(\xi, \eta) \in D^*$ :

$$p_{\xi,\eta}(x, y) = \frac{1}{4} [p_{\xi-1,\eta}(x, y) + p_{\xi+1,\eta}(x, y) + p_{\xi,\eta-1}(x, y) + p_{\xi,\eta+1}(x, y)] \quad (2)$$

this is because to get to  $(x, y)$  the particle has to walk through one of the neighbours:  $(x-1, y)$ ,  $(x+1, y)$ ,  $(x, y-1)$ ,  $(x, y+1)$ .

⇒ The expected value of the  $\nu(\xi, \eta)$  is given by equation:

$$E(\xi, \eta) = \sum_{(x,y) \in \Gamma^*} p_{\xi,\eta}(x, y) f^*(x, y) \quad (3)$$

where the summing is over all boundary points

## Laplace equation, Dirichlet boundary conditions

⇒ Now multiplying the 2 by  $f^*(x, y)$  and summing over all edge points  $(x, y)$ :

$$E(\xi, \eta) = \frac{1}{4} [E(\xi - 1, \eta) + E(\xi + 1, \eta) + E(\xi, \eta - 1) + E(\xi, \eta + 1)]$$

⇒ Putting now 1 to 3 one gets:

$$E(x, y) = f^*(x, y), \quad (\xi, \eta) \in \Gamma(D^*)$$

⇒ Now the expected value solves identical equation as our  $u^*(x, y)$  function. From this we conclude:

$$E(x, y) = u^*(x, y)$$

⇒ The algorithm:

- We put a particle in  $(x, y)$ .
- We observe it's walk up to the moment when it's on the edge  $\Gamma(D^*)$ .
- We calculate the value of  $f^*$  function in the point where the particle stops.
- Repeat the walk  $N$  times taking the average afterwards.

**Important:**

One can show the the error does not depend on the dimensions!

# Backup