## Markov Chain MC

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## Trivial example

$\Rightarrow$ Lets start with a TRIVIAL example: we want to calculate $S=A+B$. We can rewrite it in:

$$
S=p \frac{A}{p}+(1-p) \frac{B}{1-p}
$$

and one can interpret the sum as expected value of:

$$
W=\left\{\begin{array}{l}
\frac{A}{p} \text { with propability } p \\
\frac{A}{1-p} \text { with propability } 1-p
\end{array}\right.
$$

$\Rightarrow$ The algorithm:

- We generate a random variable $W$ and calculate:

$$
\hat{S}=\frac{1}{N} \sum_{i=1}^{N} W_{i}
$$

$\Rightarrow$ The $\hat{S}$ is an unbias estimator of $S$.

## Trivial example2

- Lets say we have a linear equation system:

$$
\begin{aligned}
& X=p Y+(1-p) A \\
& Y=q X+(1-q) B
\end{aligned}
$$

- We know $A, B, p, q ; X$ and $Y$ are meant to be determined.
- Algorithm:

1. We choose first element of the first equation with probability $p$ and second with probability $1-p$.
2. I we choose the second one, the outcome of this MCMC is $W=A$.
3. If we choose the first we go to second equation and choose the first element with probability $q$ and the second with $1-q$.
4. We we choose the second one, the outcome of this MCMC is $W=B$.
5. If we choose the first we go to the first equation back again.
6. We repeat the procedure.

- We can estimate the solution of this system:

$$
\hat{X}=\frac{1}{N} \sum_{i=1} W_{i} \quad \hat{\sigma_{X}}=\frac{1}{\sqrt{N-1}} \sqrt{\frac{1}{N} \sum_{i=1}^{N} W_{i}^{2}-\hat{X}^{2}}
$$

## Random walk


$\Rightarrow$ We are in the point $x$ and we walk accordingly to the following rules:

- From point $x$ we walk with probability $p$ to point $y$ or with $1-p$ to $a$.
- From point $y$ we walk with probability $q$ to point $x$ and with $1-Q$ to $b$.
- The walks ends when you end up in $a$ or $b$.
- You get a "reward" $A$ if you end up in point $a$ and $B$ if you end up in $b$.
- $X$ is expected "reward" when you start the walk from $x, Y$ when you start from $y$.
$\Rightarrow$ The algorithm above is so-called random walk on the set $\{a, x, y, b\}$
$\Rightarrow$ The described walked can solve the linear equation system that we discussed above.


## Markov Chain MC

- Consider a finite (or Countable set) possible states: $S_{1}, S_{2}, \ldots$
- The $X_{t}$ is the state of the system in the time $t$
- We are looking at discrete time steps: $1,2,3, \ldots$
- The conditional probability is defined as:

$$
P\left(X_{t}=S_{j} \mid X_{t-1}=S_{j-1}, \ldots, X_{1}=S_{1}\right)
$$

- The Markov chain is then if the probability depends only on previous step.

$$
P\left(X_{t}=S_{j} \mid X_{t-1}=S_{j-1}, \ldots, X_{1}=S_{1}\right)=P\left(X_{t}=S_{j} \mid X_{t-1}=S_{j-1}\right)
$$

- For this reason MCMC is also knows as drunk sailor walk.
- Very powerful method. Used to solve linear eq. systems, invert matrix, solve differential equations, etc.
- Also used in physics problems: Brown motions, diffusion, etc.


## Linear equations system

$\Rightarrow$ Lets start from a linear equation system:

$$
\mathbf{A} \vec{x}=\vec{b}, \quad \operatorname{det} \mathbf{A} \neq 0
$$

where $\mathbf{A}=\left(a_{i j}, i, j=1,2, \ldots, n\right.$-matrix, $\vec{b}=\left(b_{1}, b_{2}, \ldots, b_{n}\right)$-vector, $\vec{x}=\left(x_{1}, x_{2}, \ldots, x_{n}\right)-$ vector of unknowns.
$\Rightarrow$ The solution we mark as $\vec{x}^{0}=\left(x_{1}^{0}, x_{2}^{0}, \ldots, x_{n}^{0}\right)$
$\Rightarrow$ The above system can be transformed into the iterative representation:

$$
\vec{x}=\vec{a}+\mathbf{H} \vec{x}
$$

where H is a matrix, $\vec{a}$ is a vector.
$\Rightarrow$ We assume that the matrix norm:

$$
\|H\|=\max _{1 \leqslant i \leqslant n} \sum_{j=1}^{n}\left|h_{h_{i j}}\right|<1
$$

## Linear equations system

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$$

$\Rightarrow$ We can always change transform every system to the iteration form: $\mathrm{A}=\mathrm{V}-\mathrm{W}$.

$$
(\mathbf{V}-\mathbf{W}) \vec{x}=\vec{b} \quad \mapsto \quad \vec{x}=\mathbf{V}^{-1} \vec{b}+\mathbf{V}^{-1} \mathbf{W} \vec{x}
$$

## Linear equations system

$\Rightarrow$ Now we further modify the equation system:

$$
\vec{x}=\vec{a}+\mathbf{H} \vec{x} \Rightarrow(\mathbf{I}-\mathbf{H}) \vec{x}=\vec{a}
$$

where $\mathbf{I}=\delta_{i j}$ - unit matrix, $\delta_{i j}$ is the Kronecker delta.
$\Rightarrow$ What one can do is to represent the solution in terns of Neumann series:

$$
\vec{x}^{0}=(\mathbf{I}-\mathbf{H})^{-1} \vec{a}=\vec{a}+\mathbf{H} \vec{a}+\mathbf{H}^{2} \vec{a}+\mathbf{H}^{3} \vec{a}+\ldots
$$

$\Rightarrow$ So for the $i^{t h}$ component we have:

$$
\begin{aligned}
x_{i}^{0}=a_{i} & +\sum_{j=1}^{n} h_{i j} a_{j}+\sum_{j_{1}=1}^{n} \sum_{j_{2}=1}^{n} h_{i j_{1}} h_{j_{1} j_{2}} a_{j_{2}} \\
& +\ldots+\sum_{j_{1}=1}^{n} \ldots \sum_{j_{n}=1}^{n} h_{i j_{1}} \ldots h_{j_{n-1} j_{n}} a_{j_{n}}
\end{aligned}
$$

$\Rightarrow$ We will construct a probabilistic interpretation using MCMC and then we show that the expected value is equal to the above formula.

## Neumann-Ulam method

- To do so we add to our matrix an additional column of the matrix:

$$
h_{i, 0}=1-\sum_{j=1}^{n} h_{i j}>0
$$

- The system has states: $\{0,1,2 \ldots, n\}$
- State at $t$ time is denoted as $i_{t}\left(i_{t}=0,1,2, \ldots, n ; t=0,1, \ldots\right)$
- We make a random walk accordingly to to the following rules:
- At the beginning of the walk $(t=0)$ we are at $i_{0}$.
- In the $t$ moment we are in the $i_{t}$ position then in $t+1$ time stamp we move to state $i_{t+1}$ with the probability $h_{i_{t} i_{t+1}}$.
- We stop walking if we are in state 0 .
- The path $\gamma=\left(i_{0}, i_{1}, i_{2}, \ldots, i_{k}, 0\right)$ is called trajectory.
- For each trajectory we assign a number:

$$
X(\gamma)=X\left(i_{0}, i_{1}, i_{2}, \ldots, i_{k}, 0\right)=\frac{a_{i_{k}}}{h_{i_{k} 0}}
$$

## Neumann-Ulam method

$\Rightarrow$ The $X(\gamma)$ variable is a random variable from: $\left\{a_{1} / h_{1,0}, a_{2} / h_{2,0}, \ldots, a_{n} / h_{n, 0}\right\}$. The probability that $X(\gamma)=a_{j} / h_{j, 0}$ is equal to the probability that the last non zero state of the $\gamma$ trajectory is $j$.
$\Rightarrow$ The expected value of the $X(\gamma)$ trajectory if the trajectory begins from $i_{0}=s$ is:

$$
E\left\{X(\gamma) \mid i_{0}=s\right\}=\sum_{k=0}^{\infty} \sum_{\left\{\gamma_{k}\right\}} X(\gamma) P(\gamma)
$$

where $\gamma_{k}$ is a trajectory of length $k$, which starts in $i_{0}=s$ and $P(\gamma)$ is the probability of occurrence of this trajectory. $\Rightarrow$ Yes you guest it lets do Taylor expansion:

$$
E\left\{X(\gamma) \mid i_{0}=s\right\}=\sum_{\gamma_{0}} X(\gamma) P(\gamma)+\sum_{\gamma_{1}} X(\gamma) P(\gamma)+\ldots+\sum_{\gamma_{k}} X(\gamma) P(\gamma)
$$

$\Rightarrow$ Now let's examine the elements of the above series.

## Neumann-Ulam method

$\left\{\gamma_{0}\right\}$ : One trajectory: $\gamma_{0}=\left(i_{0}=s \mid 0\right), P\left(\gamma_{0}\right)=h_{s, 0}$ and $X\left(\gamma_{0}\right)=a_{s} / h_{s, 0}$. So:

$$
\sum_{\gamma_{0}} X(\gamma) P(\gamma)=\frac{a_{s}}{h_{s, 0}} h_{s, 0}=a_{s}
$$

$\left\{\gamma_{1}\right\}$ : Trajectories: $\gamma_{1}=\left(i_{0}=s, i_{1} \mid 0\right), i_{1} \neq 0, P\left(\gamma_{1}\right)=P\left(s, i_{1}, 0\right)=h_{s, i_{1}} h_{i_{1}, 0}$ and $X\left(\gamma_{1}\right)=a_{i_{1}} / h_{i_{1}, 0}$. So:

$$
\sum_{\gamma_{1}} X(\gamma) P(\gamma)=\sum_{i_{1}=1}^{n} \frac{a_{i_{1}}}{h_{i_{1}, 0}} h_{s, i_{1}} h_{i_{1}, 0}=\sum_{i=1}^{n} h_{s, i_{1}} a_{i_{1}}
$$

$\left\{\gamma_{2}\right\}$ : Trajectories: $\gamma_{2}=\left(i_{0}=s, i_{1}, i_{2} \mid 0\right), i_{1}, i_{2} \neq 0$,
$P\left(\gamma_{2}\right)=P\left(s, i_{1}, i_{2}, 0\right)=h_{s, i_{1}} h_{i_{1}, i_{2}} h_{i_{1}, 0}$ and $X\left(\gamma_{2}\right)=a_{i_{2}} / h_{i_{2}, 0}$. So:

$$
\sum_{\gamma_{2}} X(\gamma) P(\gamma)=\sum_{i_{1}=1}^{n} \sum_{i_{2}=1}^{n} \frac{a_{i_{2}}}{h_{i_{2}, 0}} h_{s, i_{1}} h_{i_{1}, i_{2}} h_{i_{2}, 0}=\sum_{i_{1}=1}^{n} \sum_{i_{2}=1}^{n} h_{s, i_{1}} h_{i_{1}, i_{2}} a_{i_{2}}
$$

$\Rightarrow$ etc...

## Neumann-Ulam method

$\Rightarrow$ After summing up:

$$
\begin{aligned}
E\left\{X(\gamma) \mid i_{0}=s\right\}=a_{s} & +\sum_{i_{1}=1}^{n} h_{s, i_{1}} a_{i_{1}}+\sum_{i_{1}=1}^{n} \sum_{i_{2}=1}^{n} h_{s, i_{1}} h_{i_{1}, i_{2}} a_{i_{2}}+\ldots \\
& +\sum_{i_{1}=1}^{n} \sum_{i_{2}=1}^{n} \ldots \sum_{i_{k}=1}^{n} h_{s, i_{1}} h_{i_{1}, i_{2}} \ldots h_{i_{k-1}, i_{k}} a_{i_{k}}+\ldots
\end{aligned}
$$

$\Rightarrow$ If you compare this expression with the Neumann series we will they are the same so:

$$
x_{i}^{0}=E\left\{X(\gamma) \mid i_{0}=i\right\}
$$

## To sum up:

We have proven that solving a linear system can be represented by an expectation value of the random variable $X(\gamma)$. The error is computed using standard deviation equation.

## Neumann-Ulam method

- For example lets try to solve this equation system:

$$
\vec{x}=\left(\begin{array}{c}
1.5 \\
-1.0 \\
0.7
\end{array}\right)+\left(\begin{array}{ccc}
0.2 & 0.3 & 0.1 \\
0.4 & 0.3 & 0.2 \\
0.3 & 0.1 & 0.1
\end{array}\right) \vec{x}
$$

- The solution is $\vec{x}_{0}=(2.154303,0.237389,1.522255)$.
- The propability matrix $h_{i j}$ has the shape:

| $i / j$ | 1 | 2 | 3 | 0 |
| :---: | :---: | :---: | :---: | :---: |
| 1 | 0.2 | 0.3 | 0.1 | 0.4 |
| 2 | 0.4 | 0.3 | 0.2 | 0.1 |
| 3 | 0.3 | 0.1 | 0.1 | 0.5 |

- An example solution:
mchrzasz-ThinkPad-W530\% ./mark.x 11000000 2.15625


## Neumann-Ulam dual method

- The problem with Neumann-Ulam method is that you need to repeat it for each of the coordinates of the $\vec{x}_{0}$ vector.
- The dual method calculates the whole $\vec{x}_{0}$ vector.
- The algorithm:
- On the indexes: $\{0,1, \ldots, n\}$ we set a probability distribution: $q_{1}, q_{2}, \ldots, q_{n}, q_{i}>0$ and $\sum_{i=1}^{n} q_{i}=1$.
- The starting point we select from $q_{i}$ distribution.
- If in $t$ time we are in $i_{t}$ state then with probability $p\left(i_{t+1} \mid i_{t}\right)=h_{i_{t+1}, i_{t}}$ in $t+1$ we will be in state $i_{t+1}$. For $i_{t+1}=0$ we define the probability: $h_{0, i_{t}}=1-\sum_{j=1}^{n} h_{j, i_{t}}$. Here we also assume that $h_{j, i_{t}}>0$.
- NOTE: there the matrix is transposed compared to previous method: $H^{T}$.
- Again we end our walk when we are at state 0 .
- For the trajectory: $\gamma=\left(i_{0}, i_{1}, \ldots, i_{k}, 0\right)$, we assign the vector:

$$
\vec{Y}(\gamma)=\frac{a_{i_{0}}}{q_{i_{0}} p\left(0 \mid i_{k}\right)} \widehat{e}_{i_{k}} \in \mathcal{R}^{n}
$$

- The solution will be : $\vec{x}^{0}=\frac{1}{N} \sum \vec{Y}(\gamma)$


## Neumann-Ulam dual method

- Let's try to solve the equation system:

$$
\vec{x}=\left(\begin{array}{c}
1.5 \\
-1.0 \\
0.7
\end{array}\right)+\left(\begin{array}{ccc}
0.2 & 0.3 & 0.1 \\
0.4 & 0.3 & 0.2 \\
0.1 & 0.1 & 0.1
\end{array}\right) \vec{x}
$$

- The solution is: $\vec{x}_{0}=(2.0,0.0,1.0)$.
- Let's put the initial probability as constant:

$$
q_{1}=q_{2}=q_{3}=\frac{1}{3}
$$

- The propability matrix $h_{i j}$ has the shape:

| $i / j$ | 1 | 2 | 3 | 4 |
| :---: | :---: | :---: | :---: | :---: |
| 1 | 0.2 | 0.4 | 0.1 | 0.3 |
| 2 | 0.3 | 0.3 | 0.1 | 0.3 |
| 3 | 0.1 | 0.2 | 0.1 | 0.6 |

- An example solution:
mchrzasz-ThinkPad-W530\% ./mark2.x 1000000 99430001806100267


## Generalization, the algorithm

$\Rightarrow$ We set the $P$ matrix in a arbitrary way.
$\Rightarrow$ If in the $t$ moment the point is in the $i_{t}$ state, then with the probability $p_{i_{t}, i_{t+1}}$ he can go to $i_{t+1}$ state.
$\Rightarrow$ We stop the walk once we reach 0 .
$\Rightarrow$ For the given trajectory we assign the value: $X\left(\gamma_{k}\right)$
$\Rightarrow$ We repeat the procedure $N$ times and take the mean and RMS.
$\Rightarrow$ We repeat this also for every of the $\vec{x}^{0}$ vector components.

## Partial differential equations, intro

$\Rightarrow$ Let's say we are want to describe a point that walks on the $\mathbb{R}$ axis:

- At the beginning $(t=0)$ the particle is at $x=0$
- If in the $t$ the particle is in the $x$ then in the time $t+1$ it walks to $x+1$ with the known probability $p$ and to the point $x-1$ with the probability $q=1-p$.
- The moves are independent.
$\Rightarrow$ So let's try to described the motion of the particle.
$\Rightarrow$ The solution is clearly a probabilistic problem. Let $\nu(x, t)$ be a probability that at time $t$ particle is in position $x$. We get the following equation:

$$
\nu(x, t+1)=p \nu(x-1, t)+q \nu(x+1, t)
$$

with the initial conditions:

$$
\nu(0,0)=1, \quad \nu(x, 0)=0 \text { if } x \neq 0
$$

$\Rightarrow$ The above functions describes the whole system (every $(t, x)$ point).

## Partial differential equations, intro

$\Rightarrow$ Now in differential equation language we would say that the particle walks in steps of $\Delta x$ in times: $k \Delta t, k=1,2,3 \ldots$ :

$$
\nu(x, t+\Delta t)=p \nu(x-\Delta x, t)+q \nu(x+\Delta x, t)
$$

$\Rightarrow$ To solve this equation we need to expand the $\nu(x, t)$ funciton in the Taylor series:

$$
\begin{aligned}
\nu(x, t)+\frac{\partial \nu(x, t)}{\partial t} \Delta t= & p \nu(x, t)-p \frac{\partial \nu(x, t)}{\partial x} \Delta x+\frac{1}{2} p \frac{\partial^{2} \nu(x, t)}{\partial x^{2}}(\Delta x)^{2} \\
& +q \nu(x, t)+q \frac{\partial \nu(x, t)}{\partial x} \Delta x+\frac{1}{2} q \frac{\partial^{2} \nu(x, t)}{\partial x^{2}}(\Delta x)^{2}
\end{aligned}
$$

$\Rightarrow$ From which we get:

$$
\frac{\partial \nu(x, t)}{\partial t} \Delta t=-(p-q) \frac{\partial \nu(x, t)}{\partial x} \Delta x+\frac{1}{2} \frac{\partial^{2} \nu(x, t)}{\partial x^{2}}(\Delta x)^{2}
$$

$\Rightarrow$ Now We divide the equation by $\Delta t$ and take the $\Delta t \rightarrow 0$ :

$$
(p-q) \frac{\Delta x}{\Delta t} \rightarrow 2 c, \quad \frac{(\Delta x)^{2}}{\Delta t} \rightarrow 2 D
$$

$\Rightarrow$ We get the Fokker-Planck equation for the diffusion with current:

$$
\frac{\partial \nu(x, t)}{\partial t}=-2 c \frac{\partial \nu(x, t)}{\partial x}+D \frac{\partial^{2} \nu(x, t)}{\partial x^{2}}
$$

## Laplace equation, Dirichlet boundary conditions

$\Rightarrow$ The aforementioned example shows the way to solve the partial differential equation using Markov Chain MC.
$\Rightarrow$ We will see how different classes of partial differential equations can be approximated with a Markov Chain MC, whose expectation value is the solution of the equation. $\Rightarrow$ The Laplace equation:

$$
\frac{\partial^{2} u}{\partial x_{1}^{2}}+\frac{\partial^{2} u}{\partial x_{2}^{2}}+\ldots+\frac{\partial^{2} u}{\partial x_{k}^{2}}=0
$$

The $u\left(x_{1}, x_{2}, \ldots, x_{k}\right)$ function that is a solution of above equation we call harmonic function. If one knows the values of the harmonic function on the edges $\Gamma(D)$ of the $D$ domain one can solve the equation.

## The Dirichlet boundary conditions:

Find the values of $u\left(x_{1}, x_{2}, \ldots, x_{k}\right)$ inside the $D$ domain knowing the values of the edge are given with a function:

$$
u\left(x_{1}, x_{2}, \ldots, x_{k}\right)=f\left(x_{1}, x_{2}, \ldots, x_{k}\right) \in \Gamma(D)
$$

$\Rightarrow$ Now I am lazy so I put $k=2$ but it's the same for all $k$ !

## Laplace equation, Dirichlet boundary conditions

$\Rightarrow$ We will put the Dirichlet boundary condition as a discrete condition:

- The domain $D$ we put a lattice with distance $h$.
- Some points we treat as inside (denoted with circles). Their form a set denoted $D^{*}$.
- The other points we consider as the boundary points and they form a set $\Gamma(D)$.

$\Rightarrow$ We express the second derivatives with the discrete form:

$$
\frac{\frac{u(x+h)-u(x)}{h}-\frac{u(x)-u(x-h)}{h}}{h}=\frac{u(x+h)-2 u(x)+u(x-h)}{h^{2}}
$$

$\Rightarrow$ Now we choose the units so $h=1$.

## Laplace equation, Dirichlet boundary conditions

## The Dirichlet condition in the discrete form:

Find the $u^{*}$ function which obeys the differential equation:

$$
U^{*}(x, y)=\frac{1}{4}\left[u^{*}(x-1, y)+u^{*}(x+1, y)+u^{*}(x, y-1)+u^{*}(x, y+1)\right]
$$

in all points $(x, y) \in D^{*}$ with the condition:

$$
u^{*}(x, y)=f^{*}(x, y), \quad(x, y) \in \Gamma\left(D^{*}\right)
$$

where $f^{*}(x, y)$ is the discrete equivalent of $f(x, y)$ function.
$\Rightarrow$ We consider a random walk over the lattice $D^{*} \cup \Gamma\left(D^{*}\right)$.

- In the $t=0$ we are in some point $\left.(\xi, \eta) \in D^{*}\right)$
- If at the $t$ the particle is in $(x, y)$ then at $t+1$ it can go with equal probability to any of the four neighbour lattices: $(x-1, y),(x+1, y),(x, y-1),(x, y+1)$.
- If the particle at some moment gets to the edge $\Gamma\left(D^{*}\right.$ then the walk is terminated.
- For the particle trajectory we assign the value of: $\nu(\xi, \eta)=f^{*}(x, y)$, where $(x, y) \in \Gamma\left(D^{*}\right)$.



## Laplace equation, Dirichlet boundary conditions

$\Rightarrow$ Let $p_{\xi, \eta}(x, y)$ be the probability of particle walk that starting in $(\xi, \eta)$ to end the walk in ( $x, y$ ).
$\Rightarrow$ The possibilities:

1. The point $(\xi, \eta) \in \Gamma\left(D^{*}\right)$. Then:

$$
p_{\xi, \eta}(x, y)= \begin{cases}1, & (x, y)=\xi, \eta)  \tag{1}\\ 0, & (x, y) \neq \xi, \eta)\end{cases}
$$

2. The point $(\xi, \eta) \in D^{*}$ :

$$
\begin{equation*}
p_{\xi, \eta}(x, y)=\frac{1}{4}\left[p_{\xi-1, \eta}(x, y)+p_{\xi+1, \eta}(x, y)+p_{\xi, \eta-1}(x, y)+p_{\xi, \eta+1}(x, y)\right] \tag{2}
\end{equation*}
$$

this is because to get to $(x, y)$ the particle has to walk through one of the neighbours: $(x-1, y),(x+1, y),(x, y-1),(x, y+1)$.
$\Rightarrow$ The expected value of the $\nu(\xi, \eta)$ is given by equation:

$$
\begin{equation*}
E(\xi, \eta)=\sum_{(x, y) \in \Gamma^{*}} p_{\xi, \eta}(x, y) f^{*}(x, y) \tag{3}
\end{equation*}
$$

where the summing is over all boundary points

## Laplace equation, Dirichlet boundary conditions

$\Rightarrow$ Now multiplying the 2 by $f^{*}(x, y)$ and summing over all edge points $(x, y)$ :

$$
E(\xi, \eta)=\frac{1}{4}[E(\xi-1, \eta)+E(\xi+1, \eta)+E(\xi, \eta-1)+E(\xi, \eta+1)]
$$

$\Rightarrow$ Putting now 1 to 3 one gets:

$$
E(x, y)=f^{*}(x, y), \quad(\xi, \eta) \in \Gamma\left(D^{*}\right)
$$

$\Rightarrow$ Now the expected value solves identical equation as our $u^{*}(x, y)$ function. From this we conclude:

$$
E(x, y)=u^{*}(x, y)
$$

$\Rightarrow$ The algorithm:

- We put a particle in $(x, y)$.
- We observe it's walk up to the moment when it's on the edge $\Gamma\left(D^{*}\right)$.
- We calculate the value of $f^{*}$ function in the point where the particle stops.
- Repeat the walk $N$ times taking the average afterwards.


## Important:

One can show the the error does not depend on the dimensions!

## Backup

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