

(Re)interpretation of Flavour Constraints

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(Re)interpreting the results of new physics searches at the LHC
CERN, December 12, 2016

Outline

- ⇒ Theoretical framework for B decays
- ⇒ $B \rightarrow K^* \ell^+ \ell^-$ observables and calculations
- ⇒ Which data do Flavour factories publish
- ⇒ New Physics searches
- ⇒ What would be the best way to exchange the information?
- ⇒ Wilson Coefficients fits with **GAMBIT**
- ⇒ Questions for discussion

Theoretical framework for B decays.

Theoretical framework for B decays

A multi-scale problem

- new physics: $\Lambda_{\text{NP}} \gtrsim \text{TeV}$
- electroweak interactions: $M_W \sim 80 \text{ GeV}$
- hadronic effects: $m_b \sim 5 \text{ GeV}$
- QCD interactions: $\Lambda_{\text{QCD}} \sim 0.2 \text{ GeV}$

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\Rightarrow Effective field theory approach:

separation between low and high energies using Operator Product Expansion

- short distance: Wilson coefficients, computed perturbatively
- long distance: local operators

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left(\sum_{i=1 \dots 10, S, P} (C_i(\mu) \mathcal{O}_i(\mu) + C'_i(\mu) \mathcal{O}'_i(\mu)) \right)$$

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New physics:

- Corrections to the Wilson coefficients: $C_i \rightarrow C_i + \Delta C_i^{\text{NP}}$
- Additional operators: $\sum C_j^{\text{NP}} \mathcal{O}_j^{\text{NP}}$

Operators

$$\mathcal{O}_1 = (\bar{s}\gamma_\mu T^a P_L c)(\bar{c}\gamma^\mu T^a P_L b)$$

$$\mathcal{O}_2 = (\bar{s}\gamma_\mu P_L c)(\bar{c}\gamma^\mu P_L b)$$

$$\mathcal{O}_3 = (\bar{s}\gamma_\mu P_L b)\sum_q(\bar{q}\gamma^\mu q)$$

$$\mathcal{O}_4 = (\bar{s}\gamma_\mu T^a P_L b)\sum_q(\bar{q}\gamma^\mu T^a q)$$

$$\mathcal{O}_5 = (\bar{s}\gamma_{\mu_1}\gamma_{\mu_2}\gamma_{\mu_3} P_L b)\sum_q(\bar{q}\gamma^{\mu_1}\gamma^{\mu_2}\gamma^{\mu_3} q)$$

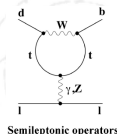
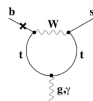
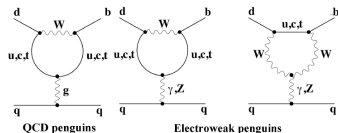
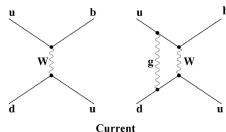
$$\mathcal{O}_6 = (\bar{s}\gamma_{\mu_1}\gamma_{\mu_2}\gamma_{\mu_3} T^a P_L b)\sum_q(\bar{q}\gamma^{\mu_1}\gamma^{\mu_2}\gamma^{\mu_3} T^a q)$$

$$\mathcal{O}_7 = \frac{e}{16\pi^2} \left[\bar{s}\sigma^{\mu\nu} (m_s P_L + m_b P_R) b \right] F_{\mu\nu}$$

$$\mathcal{O}_8 = \frac{g}{16\pi^2} \left[\bar{s}\sigma^{\mu\nu} (m_s P_L + m_b P_R) T^a b \right] G_{\mu\nu}^a$$

$$\mathcal{O}_9 = \frac{e^2}{(4\pi)^2} (\bar{s}\gamma^\mu b_L)(\bar{l}\gamma_\mu l)$$

$$\mathcal{O}_{10} = \frac{e^2}{(4\pi)^2} (\bar{s}\gamma^\mu b_L)(\bar{l}\gamma_\mu \gamma_5 l)$$



Two main steps:

- Calculating $C_i^{eff}(\mu)$ at scale $\mu \sim M_W$ by requiring matching between the effective and full theories

$$C_i^{eff}(\mu) = C_i^{(0)eff}(\mu) + \frac{\alpha_s(\mu)}{4\pi} C_i^{(1)eff}(\mu) + \dots$$

- Evolving the $C_i^{eff}(\mu)$ to scale $\mu \sim m_b$ using the RGE:

$$\mu \frac{d}{d\mu} C_i^{eff}(\mu) = C_j^{eff}(\mu) \gamma_{ji}^{eff}(\mu)$$

driven by the anomalous dimension matrix $\hat{\gamma}^{eff}(\mu)$

SM contributions to $C_i(\mu_b)$ are known to NNLO QCD and NLO EW/QED

Hadronic quantities

To compute the amplitudes:

$$\mathcal{A}(A \rightarrow B) = \langle B | \mathcal{H}_{\text{eff}} | A \rangle = \frac{G_F}{\sqrt{2}} \sum_i \lambda_i C_i(\mu) \langle B | \mathcal{O}_i | A \rangle(\mu)$$

$\langle B | \mathcal{O}_i | A \rangle$: hadronic matrix element

How to compute matrix elements?

→ Model building, Lattice simulations, Light flavour symmetries, Heavy flavour symmetries, ...

→ Describe hadronic matrix elements in terms of **hadronic quantities**

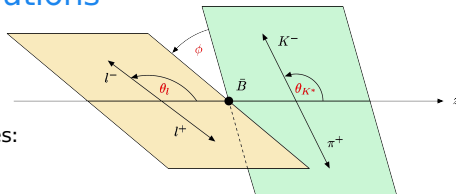
Two types of hadronic quantities:

- **Decay constants**: Probability amplitude of hadronising quark pair into a given hadron
- **Form factors**: Transition from a meson to another through flavour change

$B \rightarrow K^* \ell^+ \ell^-$ – Angular distributions

Angular distributions

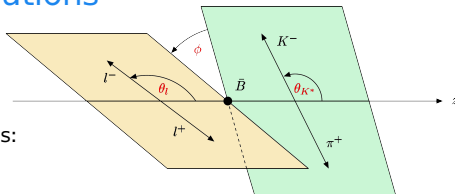
The full angular distribution of the decay $\bar{B}^0 \rightarrow \bar{K}^{*0} \ell^+ \ell^-$ ($\bar{K}^{*0} \rightarrow K^- \pi^+$) is completely described by four independent kinematic variables: q^2 (dilepton invariant mass squared), θ_ℓ , θ_{K^*} , ϕ



$B \rightarrow K^* \ell^+ \ell^-$ – Angular distributions

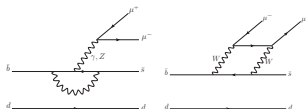
Angular distributions

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Main operators:

$$\mathcal{O}_9 = \frac{e^2}{(4\pi)^2} (\bar{s} \gamma^\mu b_L) (\bar{\ell} \gamma_\mu \ell), \quad \mathcal{O}_{10} = \frac{e^2}{(4\pi)^2} (\bar{s} \gamma^\mu b_L) (\bar{\ell} \gamma_\mu \gamma_5 \ell)$$



F. Kruger et al., Phys. Rev. D 61 (2000) 114028;

W. Altmannshofer et al., JHEP 0901 (2009) 019; U. Egede et al., JHEP 1010 (2010) 056

Differential decay distribution:

$$\frac{d^4 \Gamma}{dq^2 d \cos \theta_\ell d \cos \theta_V d \phi} = \frac{9}{32\pi} J(q^2, \theta_\ell, \theta_V, \phi)$$

$$J(q^2, \theta_\ell, \theta_V, \phi) = \sum_i J_i(q^2) f_i(\theta_\ell, \theta_V, \phi)$$

↘ angular coefficients J_{1-9}

↘ functions of the transversity amplitudes $A_0, A_{\parallel}, A_{\perp}, A_t,$

and A_S , Transversity amplitudes: functions of Wilson coefficients and form factors

$B \rightarrow K^* \ell^+ \ell^-$ – Amplitudes

A closer look to the Effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} = \mathcal{H}_{\text{eff}}^{\text{had}} + \mathcal{H}_{\text{eff}}^{\text{sl}}$$
$$\mathcal{H}_{\text{eff}}^{\text{sl}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left[\sum_{i=7,9,10} C_i^{(\prime)} O_i^{(\prime)} \right]$$

$\langle \bar{K}^* | \mathcal{H}_{\text{eff}}^{\text{sl}} | \bar{B} \rangle$: $B \rightarrow K^*$ form factors $V, A_{0,1,2}, T_{1,2,3}$

Transversity amplitudes:

$$A_{\perp}^{L,R} \simeq N_{\perp} \left\{ (C_9^+ \mp C_{10}^+) \frac{V(q^2)}{m_B + m_{K^*}} + \frac{2m_b}{q^2} C_7^+ T_1(q^2) \right\}$$

$$A_{\parallel}^{L,R} \simeq N_{\parallel} \left\{ (C_9^- \mp C_{10}^-) \frac{A_1(q^2)}{m_B - m_{K^*}} + \frac{2m_b}{q^2} C_7^- T_2(q^2) \right\}$$

$$A_0^{L,R} \simeq N_0 \left\{ (C_9^- \mp C_{10}^-) [(\dots) A_1(q^2) + (\dots) A_2(q^2)] \right. \\ \left. + 2m_b C_7^- [(\dots) T_2(q^2) + (\dots) T_3(q^2)] \right\}$$

$$A_S = N_S (C_S - C'_S) A_0(q^2)$$

$$(C_i^{\pm} \equiv C_i \pm C'_i)$$

$B \rightarrow K^* \ell^+ \ell^-$ – Amplitudes

A closer look to the Effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} = \mathcal{H}_{\text{eff}}^{\text{had}} + \mathcal{H}_{\text{eff}}^{\text{sl}}$$

$$\mathcal{H}_{\text{eff}}^{\text{had}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \left[\sum_{i=1 \dots 6} C_i O_i + C_8 O_8 \right]$$

$$\begin{aligned} \mathcal{A}_\lambda^{(\text{had})} &= -i \frac{e^2}{q^2} \int d^4 x e^{-iq \cdot x} \langle \ell^+ \ell^- | j_\mu^{\text{em, lept}}(x) | 0 \rangle \\ &\quad \times \int d^4 y e^{iq \cdot y} \langle \bar{K}_\lambda^* | T \{ j^{\text{em, had}, \mu}(y) \mathcal{H}_{\text{eff}}^{\text{had}}(0) \} | \bar{B} \rangle \\ &\equiv \frac{e^2}{q^2} \epsilon_\mu L_V^\mu \left[\underbrace{\text{LO in } \mathcal{O}\left(\frac{\Lambda}{m_b}, \frac{\Lambda}{E_{K^*}}\right)}_{\text{Non-Fact., QCDf}} + \underbrace{h_\lambda(q^2)}_{\text{power corrections}} \right] \end{aligned}$$

$B \rightarrow K^* \ell^+ \ell^-$ – Amplitudes

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$B \rightarrow K^* \ell^+ \ell^-$ – Amplitudes

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The observed deviations from the SM can be explained with 20-50% non-factorisable power corrections at the observable level (Ciuchini et al., 1512.07157)

This corresponds to more than 150% error at the amplitude level for the critical bins!

$B \rightarrow K^* \mu^+ \mu^-$ – Optimized observables

$$\begin{aligned}\langle P_1 \rangle_{\text{bin}} &= \frac{1}{2} \frac{\int_{\text{bin}} dq^2 [J_3 + \bar{J}_3]}{\int_{\text{bin}} dq^2 [J_{2s} + \bar{J}_{2s}]} & \langle P_2 \rangle_{\text{bin}} &= \frac{1}{8} \frac{\int_{\text{bin}} dq^2 [J_{6s} + \bar{J}_{6s}]}{\int_{\text{bin}} dq^2 [J_{2s} + \bar{J}_{2s}]} \\ \langle P'_4 \rangle_{\text{bin}} &= \frac{1}{\mathcal{N}'_{\text{bin}}} \int_{\text{bin}} dq^2 [J_4 + \bar{J}_4] & \langle P'_5 \rangle_{\text{bin}} &= \frac{1}{2\mathcal{N}'_{\text{bin}}} \int_{\text{bin}} dq^2 [J_5 + \bar{J}_5] \\ \langle P'_6 \rangle_{\text{bin}} &= \frac{-1}{2\mathcal{N}'_{\text{bin}}} \int_{\text{bin}} dq^2 [J_7 + \bar{J}_7] & \langle P'_8 \rangle_{\text{bin}} &= \frac{-1}{\mathcal{N}'_{\text{bin}}} \int_{\text{bin}} dq^2 [J_8 + \bar{J}_8]\end{aligned}$$

with

$$\mathcal{N}'_{\text{bin}} = \sqrt{-\int_{\text{bin}} dq^2 [J_{2s} + \bar{J}_{2s}] \int_{\text{bin}} dq^2 [J_{2c} + \bar{J}_{2c}]}$$

+ CP violating clean observables and other combinations

U. Egede et al., JHEP 0811 (2008) 032, JHEP 1010 (2010) 056

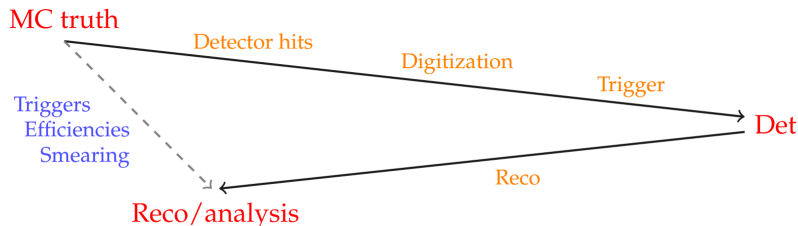
J. Matias et al., JHEP 1204 (2012) 104

S. Descotes-Genon et al., JHEP 1305 (2013) 137

Flavour measurements

Detector effects 1/2

⇒ In Flavour factories because we usually measure the properties of a B meson decay we can provide the measurements that are corrected for the detector effects!



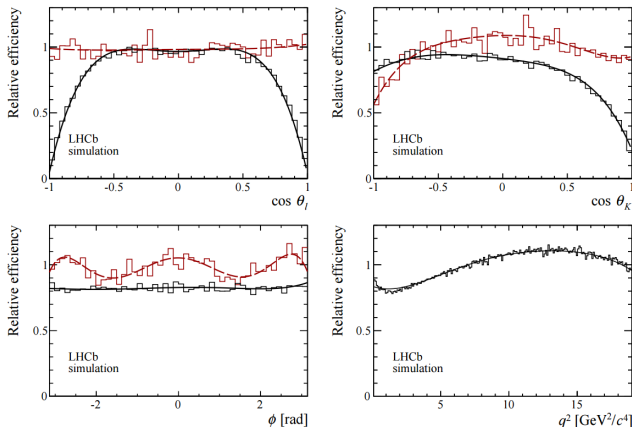
⇒ The differences that "Reco recovery" doesn't recover are recovered at the analysis stage.

⇒ Some imperfections (usually small), are assigned as systematics!

Thanks to Andy Buckley for the plot.

Detector effects 2/2

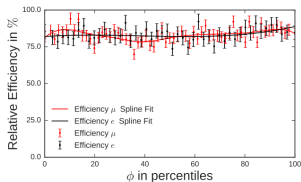
⇒ For example: measurement of angular coefficients of $B \rightarrow K^* \mu \mu$, [arXiv::1512.04442](https://arxiv.org/abs/1512.04442),
[arXiv::1604.04042](https://arxiv.org/abs/1604.04042)



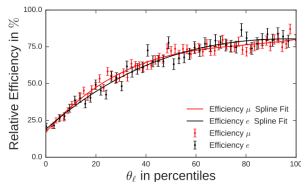
⇒ In Flavour physics we have ways to ensure we control our detector effects.

Detector effects 2/2

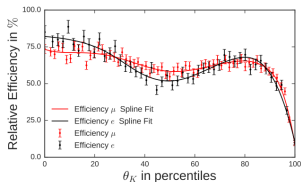
⇒ For example: measurement of angular coefficients of $B \rightarrow K^* \mu \mu$, [arXiv::1512.04442](https://arxiv.org/abs/1512.04442),
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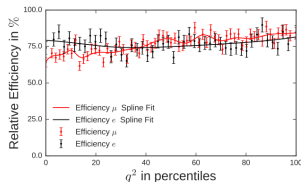
(a) Efficiency in ϕ



(b) Efficiency in θ_ℓ



(c) Efficiency in θ_K

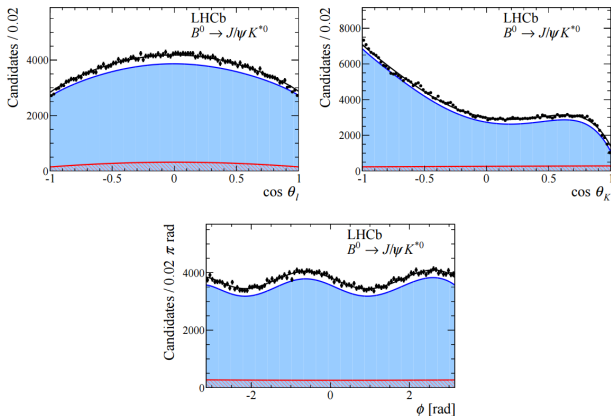


(d) Efficiency in q^2

⇒ In Flavour physics we have ways to ensure we control our detector effects.

Detector effects 2/2

⇒ For example: measurement of angular coefficients of $B \rightarrow K^* \mu \mu$, [arXiv::1512.04442](https://arxiv.org/abs/1512.04442),
[arXiv::1604.04042](https://arxiv.org/abs/1604.04042)



⇒ In Flavour physics we have ways to ensure we control our detector effects.

CERN document server



Angular analysis of the $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decay using $3. \text{fb}^{-1}$ of integrated luminosity - Aaij, Roel et al - arXiv:1512.04442

Information Discussion (5) Files

Main file(s):

- JHEP02(2016)104
version 1 JHEP02(2016)104.pdf [6.46 MiB] 23 Mar 2016, 14:44 Springer Open Access article
- arXiv:1512.04442
version 2 arXiv:1512.04442.pdf [3.43 MiB] 09 Mar 2016, 00:56
(see previous)

Additional file(s):

- LHCb-PAPER-2015-051-figures
version 1 LHCb-PAPER-2015-051-figures.zip [6.5 MiB] 11 Jan 2016, 15:10 Related data file(s)
- LHCb-PAPER-2015-051-supplementary-updated
version 1 LHCb-PAPER-2015-051-supplementary-updated.zip [33.73 MiB] 17 May 2016, 17:27 Related supplementary data file(s) updated

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- ⇒ Supplementary material not included in the paper (usually material that did not fit paper due to space constraints)

Angular analysis of the $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ decay using 3 fb^{-1} of integrated luminosity

[to restricted-access page]

INFORMATION

LHCb-PAPER-2015-051

PH-EP-2015-314

ARXIV:1512.04442 [PDF]

(SUBMITTED ON 14 DEC 2015)

JHEP 02 (2016) 104

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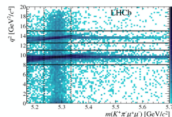
Abstract

An angular analysis of the $B^0 \rightarrow K^{*0} (\rightarrow K^+ \pi^-) \mu^+ \mu^-$ decay is presented. The dataset corresponds to an integrated luminosity of 3.0 fb^{-1} of pp collision data collected at the LHCb experiment. The complete angular information from the decay is used to determine CP -averaged observables and CP asymmetries, taking account of possible contamination from decays with the $K^+ \pi^-$ system in a 5-wave configuration. The angular observables and their correlations are reported in bins of q^2 , the invariant mass squared of the dimuon system. The observables are determined both from an unbinned maximum likelihood fit and by using the principal moments of the angular distribution. In addition, by fitting for q^2 -dependent decay amplitudes in the region $1.1 < q^2 < 6.0 \text{ GeV}^2/c^4$, the zero-crossing points of several angular observables are computed. A global fit is performed to the complete set of CP -averaged observables obtained from the maximum likelihood fit. This fit indicates differences with predictions based on the Standard Model at the level of 3.4 standard deviations. These differences could be explained by contributions from physics beyond the Standard Model, or by an unexpectedly large hadronic effect that is not accounted for in the Standard Model predictions.

Figures and captions

Invariant mass of the $K^+ \pi^- \mu^+ \mu^-$ system versus q^2 . The decay $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ is clearly visible inside the dashed vertical lines. The horizontal lines denote the charmonium regions, where the tree-level decays $B^0 \rightarrow J/\psi K^{*0}$ and $B^0 \rightarrow \psi(2S) K^{*0}$ dominate. These candidates are excluded from the analysis.

Fig1.pdf [30 KIB]
 H1Def.png [700 KIB]
 Thumbnail [205 KIB]



- ⇒ Figure on CDS and LHCb publications page available in many formats: .pdf, .eps, .png, ROOT_.C
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Angular analysis of the $B^0 \rightarrow K^+ \mu^+ \mu^-$ decay using 3 fb^{-1} of integrated luminosity

The LHCb collaboration

Aaij, Roel, Abellán Beteta, Carlos, Adro, Bernardo, Adrohn, Marco, Aflander, Anthony, Ajlouni, Ziad, Alkan Simon, Albrecht, Johannes, Altmann, Federico, Alexander, Michael

JHEP 1602 (2016) 104, 2016

http://dx.doi.org/10.17132/hepdata.74247

DOI INSPIRE Record HepData

Abstract (data abstract)
 CERN-LHC: An angular analysis of the $B^0 \rightarrow K^+ (\rightarrow K^+ \pi^-) \mu^+ \mu^-$ decay is presented. The dataset corresponds to an integrated luminosity of 3.0 fb^{-1} of pp collision data collected at the LHCb experiment. The complete angular information from the decay is used to determine CP-averaged observables and CP asymmetries, taking account of possible contamination from decays with the $K^* \pi^-$ system in an S-wave configuration. The angular observables and their correlations are reported in bins of q^2 , the invariant mass squared of the dimuon system. The observables are determined both from an unbinned maximum likelihood fit and by using the principal moments of the angular distribution. In addition, by fitting for q^2 -dependent decay amplitudes in the region $1.1 < q^2 < 6.0 \text{ GeV}^2/c^4$, the zero-crossing points of several angular observables are computed. A global fit is performed to the complete set of CP-averaged observables obtained from the maximum

Table 1
 Data from Appendix A, Table 3
 10.17132/hepdata.74247.v1.v3
 CP-averaged angular observables evaluated by the unbinned maximum likelihood fit.

Table 2
 Data from Appendix A, Table 4
 10.17132/hepdata.74247.v1.v3
 CP-asymmetric angular observables evaluated by the unbinned maximum likelihood fit. The first uncertainties are statistical and the second systematic.

Table 3
 Data from Appendix A, Table 3
 10.17132/hepdata.74247.v1.v3
 CP-averaged angular observables evaluated by the unbinned maximum likelihood fit. The first uncertainties are statistical and the second systematic.

Table 4
 Data from Appendix A, Table 4
 10.17132/hepdata.74247.v1.v4
 Optimized angular observables evaluated by the unbinned maximum likelihood fit. The first uncertainties are statistical and the second

RE	PP → BO + K*(B92) + K+ P1- → MU+ MU- + X					
SQRT(S)	7000.0 GeV					
SQRT(T)	8000.0 GeV					
q ² [GeV ²]	F _L	S ₀	S ₁	S ₂	A _{FB}	S _T
0.1 - 0.98	0.263 +0.022 stat +0.017 sys	-0.036 +0.063 stat +0.008 sys	0.082 +0.022 stat +0.009 sys	0.17 +0.022 stat +0.018 sys	-0.003 +0.009 stat +0.009 sys	0.015 +0.009 stat +0.004 sys
1.1 - 2.5	0.66 +0.022 stat +0.022 sys	-0.077 +0.035 stat +0.005 sys	-0.077 +0.035 stat +0.005 sys	0.137 +0.022 stat +0.009 sys	-0.191 +0.022 stat +0.012 sys	-0.219 +0.022 stat +0.004 sys
2.5 - 4	0.876 +0.022 stat +0.017 sys	0.035 +0.037 stat +0.007 sys	-0.234 +0.037 stat +0.004 sys	-0.022 +0.022 stat +0.011 sys	-0.118 +0.022 stat +0.007 sys	0.068 +0.022 stat +0.003 sys
4 - 6	0.611 +0.022 stat +0.017 sys	0.035 +0.037 stat +0.007 sys	-0.219 +0.037 stat +0.004 sys	-0.146 +0.022 stat +0.011 sys	0.025 +0.022 stat +0.004 sys	-0.016 +0.022 stat +0.004 sys
6 - 8	0.579 +0.022 stat +0.017 sys	-0.042 +0.035 stat +0.011 sys	-0.296 +0.035 stat +0.011 sys	-0.249 +0.022 stat +0.007 sys	0.152 +0.022 stat +0.008 sys	-0.047 +0.022 stat +0.003 sys
11 - 12.5	0.493 +0.022 stat +0.017 sys	-0.189 +0.035 stat +0.011 sys	-0.283 +0.035 stat +0.011 sys	-0.327 +0.022 stat +0.007 sys	0.318 +0.022 stat +0.008 sys	-0.141 +0.022 stat +0.003 sys

Visualize

Sum errors Log Scale (X)

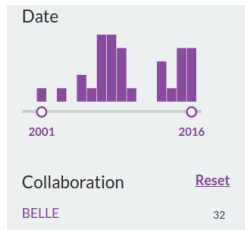
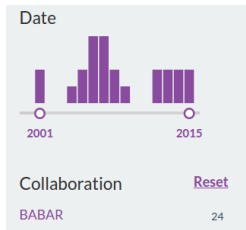
Select variables or hide different error bars by clicking on them.

Variables

F_L Summed error

Unification of format

⇒ More and more papers from Flavour community are appearing on HepData.



This is not the end of the story!!

⇒ Even if experimentalist publish a number there is always a chance that the data might be misinterpreted by theorists.

This is not the end of the story!!

- ⇒ Even if experimentalist publish a number there is always a chance that the data might be misinterpreted by theorists.
- ⇒ Many times the error gets symmetrized, the correlation neglected, or worse...

Publish likelihood?

- ⇒ The proposal that I would like to make for discussion is that HepData portal (or similar) would have a possibility that experiments could publish the whole multidim. likelihood function.
- ⇒ In this way we ensure that the function will be used as the experiment intended to.

Global fits

GAMBIT: a *second-generation* global fit code

GAMBIT: The **G**lobal **A**nd **M**odular **B**SM **I**nference **T**ool

Overriding principles of GAMBIT: flexibility and modularity

- General enough to allow fast definition of new datasets and theoretical models
- Plug and play scanning, physics and likelihood packages
- Extensive model database – not just small modifications to constrained MSSM (NUHM, etc), and not just SUSY!
- Extensive observable/data libraries (likelihood modules)
- Many statistical options – Bayesian/frequentist, likelihood definitions, scanning algorithms
- A smart and *fast* LHC likelihood calculator
- Massively parallel
- Full open-source code release soon!
- Hear more in Anders Kvellestad tmr!

The GAMBIT Collaboration

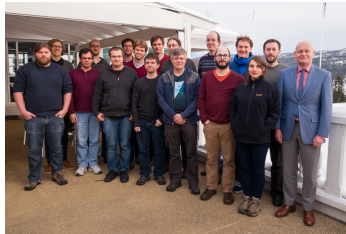
30 Members, 16 institutions, 10 countries, 11
Experiments, 4 major theory codes



ATLAS	A. Buckley, C. Rogan, M. White,
Flavour exp.	F. Bernlochner, M. Chrzaszcz, P. Jackson, N. Serra
Fermi-LAT	J. Conrad, J. Edsjö, G. Martinez P. Scott
CTA	C. Balázs, T. Bringmann, J. Conrad, M. White
HESS	J. Conrad
IceCube	J. Edsjö, P. Scott
AMS-02	A. Putze
CDMS, DM-ICE	L. Hsu
XENON/DARWIN	J. Conrad
Theory	P. Athron, C. Balázs, T. Bringmann, J. Cornell, L. Dal, J. Edsjö, B. Farmer, A. Krislock, A. Kvellestad, M. Pato, F. Mahmoudi, A. Raklev, P. Scott, C. Weniger, M. White

+recently joined: T. Gonzales, J. McKay, R. Ruiz, R. Trotta

-recently retired: A. Saavedra, C. Savage



Global Analysis with Gambit

- Wilson coefficients and $b \rightarrow s\ell^+\ell^-$ observables implemented in **SuperIso**
- **SuperIso**: public code for calculating flavour physics observables

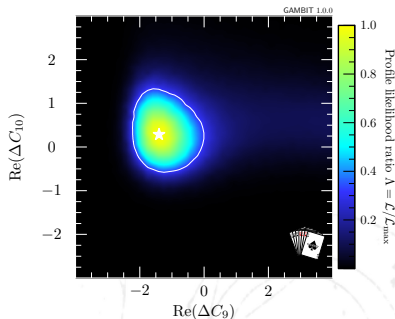
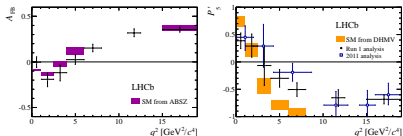
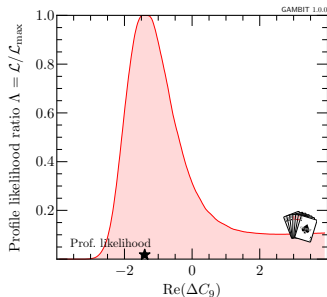
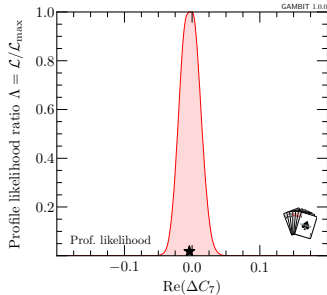
Mahmoudi, CPC 178 (2008) 745; CPC 180 (2009) 1579, CPC 180 (2009) 1718
available from <http://superiso.in2p3.fr/>

- **SuperIso** interfaced into **GAMBIT** through the flavour physics module **FlavBit**

Web page: <http://gambit.hepforge.org/>

- **FlavBit** determines the likelihoods by comparing the theoretical evaluations and the experimental results taking into account the experimental and theoretical correlations.
- In this study we used:
 - $B \rightarrow K^* \mu\mu$ with all the q^2 bins and correlations matrices from HepData!
 - $B_{s/d} \rightarrow \mu\mu$
 - $b \rightarrow s\gamma$

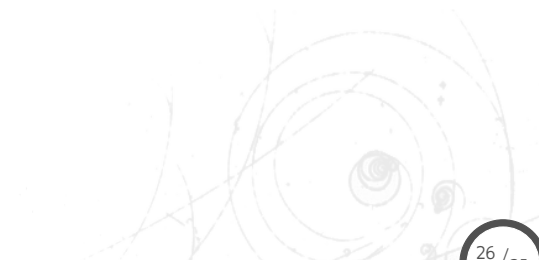
Global Analysis with Gambit - Results



- ⇒ Tension if ΔC_9 observed!
- ⇒ Other coefficients within SM predictions.
- ⇒ C_{10} still has a big uncertainty.

Conclusions

- ⇒ Flavour physics is a powerful tool to constrain NP models!
- ⇒ Measurements are becoming more complex!
- ⇒ Ability to publish the full multidim. likelihoods soon will be needed!
- ⇒ **GAMBIT** is the new player for fitting Flavour observables and will be made public soon.
- ⇒ $3-4\sigma$ deviations are present and Run2 data should clear the picture where it's NP or not.



Numerical approach

