

Searches for heavy neutral leptons at the Future Circular Colliders



Marcin Chrzaszcz
mchrzasz@cern.ch



on behalf of the FCC collaboration,
CERN and

Institute of Nuclear Physics, Polish Academy of Science

15th Rencontres du Vietnam,
Quy Nhon, 4-10 August 2019

Particles of SM

SM

| | | | |
|----------|------------------------------|-----------------------------|------------------------------|
| mass → | 2.4 MeV | 1.27 GeV | 171.2 GeV |
| charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ |
| name → | u | c | t |
| | Left up Right | Left charm Right | Left top Right |
| Quarks | 4.8 MeV | 104 MeV | 4.2 GeV |
| | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ |
| | d | s | b |
| | Left down Right | Left strange Right | Left bottom Right |
| Leptons | 0 eV | 0 eV | 0 eV |
| | 0 | 0 | 0 |
| | ν_e | ν_μ | ν_τ |
| | Left electron neutrino Right | Left muon neutrino Right | Left tau neutrino Right |
| | 0.511 MeV | 105.7 MeV | 1.777 GeV |
| | -1 | -1 | -1 |
| | e | μ | τ |
| | Left electron Right | Left muon Right | Left tau Right |

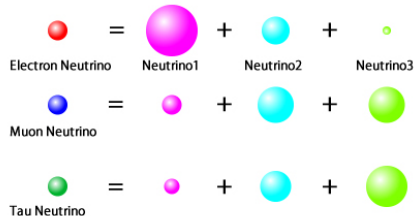
Particles of SM

| SM | | | | | | |
|----------|---------------------------|-----------------------------|------------------------------|--|--|--|
| mass → | 2.4 MeV | 1.27 GeV | 171.2 GeV | | | |
| charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | | | |
| name → | u | c | t | | | |
| | Left up Right | Left charm Right | Left top Right | | | |
| | 4.8 MeV | 104 MeV | 4.2 GeV | | | |
| | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | | | |
| Quarks | d | s | b | | | |
| | Left down Right | Left strange Right | Left bottom Right | | | |
| | 0 eV | 0 eV | 0 eV | | | |
| | 0 | 0 | 0 | | | |
| | ν_e | ν_μ | ν_τ | | | |
| | electron neutrino | muon neutrino | tau neutrino | | | |
| | 0.511 MeV | 105.7 MeV | 1.777 GeV | | | |
| | -1 | -1 | -1 | | | |
| Leptons | e | μ | τ | | | |
| | Left electron Right | Left muon Right | Left tau Right | | | |

⇒ In the SM neutrinos do not appear in the right-handed state.

⇒ By construction neutrinos are massless.

⇒ Neutrino oscillations are evidence for physics beyond the SM!



Extending the SM

⇒ A lazy person solution is to add the right-handed neutrinos in:

| SM | | | | nuMSM | | | |
|----------|------------------------------------------------|----------------------------------------------|----------------------------------------------|------------------------------------------------|----------------------------------------------|----------------------------------------------|---------------------------------------------|
| mass → | 2.4 MeV | 1.27 GeV | 171.2 GeV | 2.4 MeV | 1.27 GeV | 171.2 GeV | |
| charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | |
| name → | u up | c charm | t top | u up | c charm | t top | |
| | Left Right | Left Right | Left Right | Left Right | Left Right | Left Right | |
| | 4.8 MeV | 104 MeV | 4.2 GeV | 4.8 MeV | 104 MeV | 4.2 GeV | |
| | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | |
| Quarks | d down | s strange | b bottom | d down | s strange | b bottom | |
| | Left Right | Left Right | Left Right | Left Right | Left Right | Left Right | |
| | 0 eV | 0 eV | 0 eV | <0.0001 eV | ~10 keV | ~0.01 eV | ~GeV |
| | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| | ν_e electron neutrino | ν_μ muon neutrino | ν_τ tau neutrino | ν_e electron neutrino | ν_μ muon neutrino | ν_τ tau neutrino | N_1 sterile neutrino |
| | Left Right | Left Right | Left Right | Left Right | Left Right | Left Right | Left Right |
| | 0.511 MeV | 105.7 MeV | 1.777 GeV | 0.511 MeV | 105.7 MeV | 1.777 GeV | |
| | -1 | -1 | -1 | -1 | -1 | -1 | |
| Leptons | e electron | μ muon | τ tau | e electron | μ muon | τ tau | |
| | Left Right | Left Right | Left Right | Left Right | Left Right | Left Right | |

Extending the SM

⇒ A lazy person solution is to add the right-handed neutrinos in:

| SM | | | | | | nuMSM | | | | | | | | | |
|----------|-------------------|----------------|----------------|----------|-------------------|----------------|----------------|------------------|---------------|---------------|---------------|----------|---------------|---------------|---------------|
| mass → | 2.4 MeV | 1.27 GeV | 171.2 GeV | mass → | 2.4 MeV | 1.27 GeV | 171.2 GeV | charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ |
| charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | charge → | $\frac{2}{3}$ | $\frac{2}{3}$ | $\frac{2}{3}$ | name → | Left | Right | Left | Right | Left | Right | Left |
| name → | u | c | t | name → | u | c | t | Quarks | Left | Right | Left | Right | Left | Right | Left |
| | up | charm | top | | up | charm | top | | Right | Right | Right | Right | Right | Right | Right |
| | 4.8 MeV | 104 MeV | 4.2 GeV | | 4.8 MeV | 104 MeV | 4.2 GeV | | Left | Right | Left | Right | Left | Right | Left |
| | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | | $-\frac{1}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{3}$ | | Left | Right | Left | Right | Left | Right | Left |
| Quarks | d | s | b | | d | s | b | | Left | Right | Left | Right | Left | Right | Left |
| | down | strange | bottom | | down | strange | bottom | | Left | Right | Left | Right | Left | Right | Left |
| | 0 eV | 0 eV | 0 eV | | <0.0001 eV | ~ 10 keV | ~ 0.01 eV | \sim GeV | | Left | Right | Left | Right | Left | Right |
| | ν_e | ν_μ | ν_τ | | ν_e | ν_μ | ν_τ | N_1 | | Left | Right | Left | Right | Left | Right |
| | electron neutrino | muon neutrino | tau neutrino | | electron neutrino | muon neutrino | tau neutrino | sterile neutrino | | Left | Right | Left | Right | Left | Right |
| | 0.511 MeV | 105.7 MeV | 1.777 GeV | | 0.511 MeV | 105.7 MeV | 1.777 GeV | | Left | Right | Left | Right | Left | Right | Left |
| | -1 | -1 | -1 | | -1 | -1 | -1 | | Left | Right | Left | Right | Left | Right | Left |
| Leptons | e | μ | τ | Leptons | e | μ | τ | | Left | Right | Left | Right | Left | Right | Left |
| | electron | muon | tau | | electron | muon | tau | | Left | Right | Left | Right | Left | Right | Left |

⇒ But where are they?

arXiv::hep-ph/0605047, M.Shaposhnikov

Seesaw mechanism

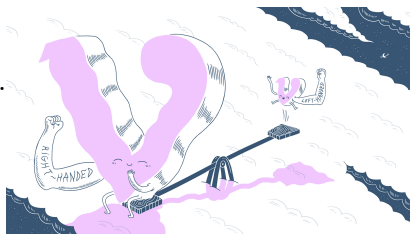
$$\mathcal{L} = \mathcal{L}_{SM} + \bar{\ell}_L F \nu_R \epsilon \Phi^* - \frac{1}{2} \bar{\nu}_R^c M_M \nu_R + \text{H.c.}$$

⇒ After the EWSB:

$$\frac{1}{2} (\bar{\nu}_L \bar{\nu}_R^c) \mathcal{M} (\nu_L^c \nu_R)^T$$

⇒ In the vanilla seesaw:

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_M \end{pmatrix} \quad \Rightarrow \lambda_+ \sim M_D, \quad \lambda_- \sim -\frac{M_M^2}{M_D}$$



Seesaw mechanism

$$\mathcal{L} = \mathcal{L}_{SM} + \bar{\ell}_L F \nu_R \epsilon \Phi^* - \frac{1}{2} \bar{\nu}_R^c M_M \nu_R + \text{H.c.}$$

⇒ After the EWSB:

$$\frac{1}{2} (\bar{\nu}_L \bar{\nu}_R^c) \mathcal{M} (\nu_L^c \nu_R)^T$$

⇒ In the vanilla seesaw:

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_M \end{pmatrix} \Rightarrow \lambda_+ \sim M_D, \quad \lambda_- \sim -\frac{M_M^2}{M_D}$$

⇒ In reality:

$$\mathcal{M} = \begin{pmatrix} \delta m_\nu^{1loop} & M_D \\ M_D^T & M_M + \delta M_N^{1loop} \end{pmatrix}$$



Seesaw mechanism

⇒ Diagonalization matrix:

$$\mathcal{U} = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta^\dagger) & \cos(\theta^\dagger) \end{pmatrix} \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix}$$

$$\mathcal{U}^\dagger \mathcal{M} \mathcal{U}^* = \begin{pmatrix} m_\nu^{\text{diag}} & \\ & M_N^{\text{diag}} \end{pmatrix}$$

with

$$M_N^{\text{diag}} = U_N^T M_N U_N = \text{diag}(M_1, M_2, M_3)$$

$$m_\nu^{\text{diag}} = U_\nu^\dagger m_\nu U_\nu^* = \text{diag}(m_1, m_2, m_3).$$

⇒ For small mixings:

$$\mathcal{U} = \left[\begin{pmatrix} \mathbb{I} - \frac{1}{2}\theta\theta^\dagger & \theta \\ -\theta^\dagger & \mathbb{I} - \frac{1}{2}\theta^\dagger\theta \end{pmatrix} + \mathcal{O}(\theta^3) \right] \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix},$$

Correction to SM processes

⇒ Charge currents:

$$j_{\mu}^{+} = \frac{g}{2} \theta_{\alpha} e \bar{l} l_{\alpha} \gamma_{\mu} N$$

⇒ Neutral currents:

$$j_{\mu}^0 = \nu_{\alpha} \gamma_{\mu} \theta_{\alpha} N$$

⇒ The Yukawa couplings:

$$\mathcal{L}_{Yukawa} = \mathcal{L}_{Yukawa}^{SM} \theta_{\alpha}$$

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

Indirect

- EW precision observables
- LFV, LNV
- Neutrinoless double beta decay
- Big Bang Nucleosynthesis

Current status

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

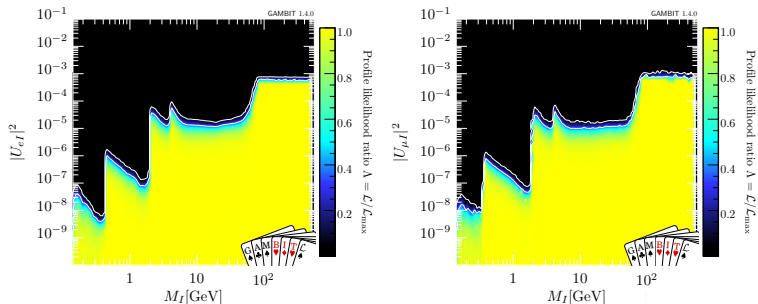
Indirect

- EW precision observables
- LFV, LNV
- Neutrinoless double beta decay
- Big Bang Nucleosynthesis

Direct

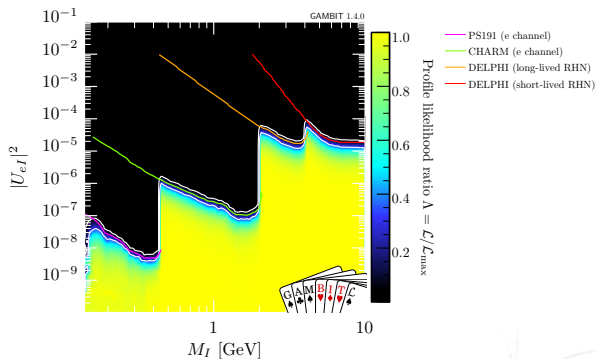
- Beam dump searches
- Collider searches (LEP, LHC, etc.)

Current status



⇒ M.C., M. Drewers, T. Gonzalo, J. Harz, S. Krishnamurthy, C. Weniger, arXiv::1908:XYZZ

Where the constrains come from?

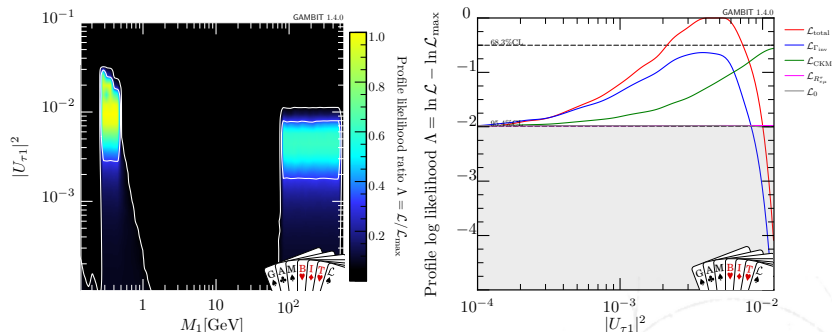


⇒ The direct searches are the strongest constraints where production cross sections are the largest.

⇒ arXiv::1908:XYZZ

Indirect constraints

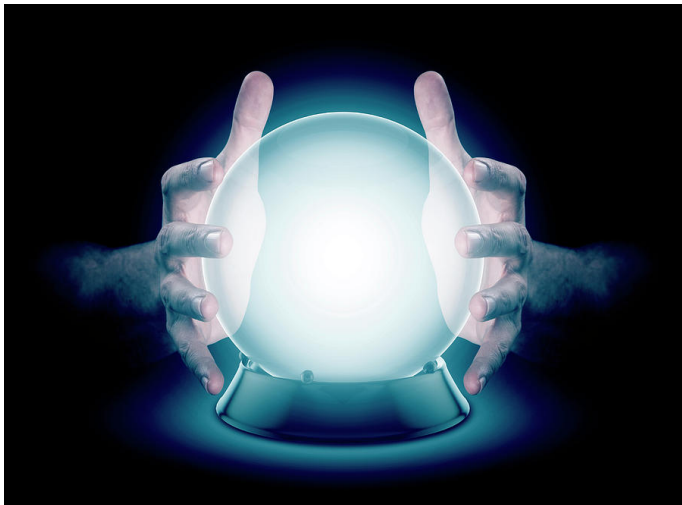
⇒ The indirect searches show power for high couplings:



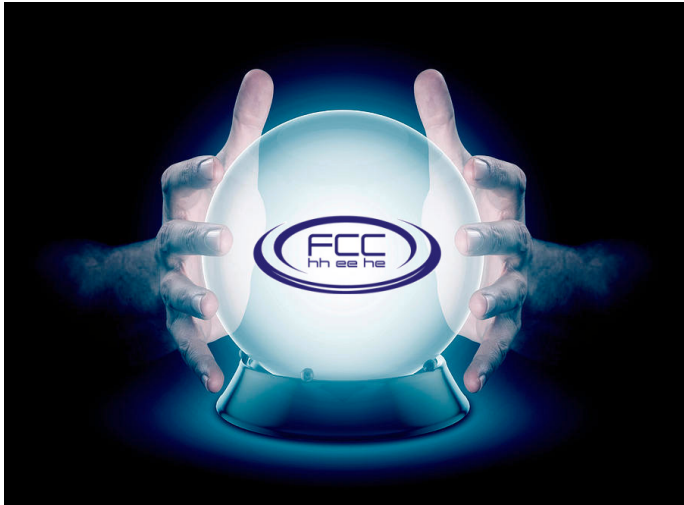
⇒ Small excess is visible. Consistent with the fluctuation.

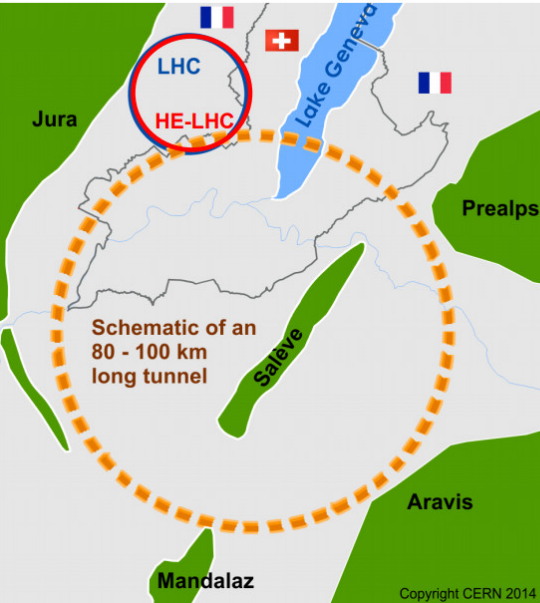
⇒ arXiv::1908:XYZZ

What can happen in the future?



What can happen in the future?





International FCC collaboration with CERN as host lab to study:

- ~100 km tunnel infrastructure in Geneva area, linked to CERN
- e^+e^- collider (*FCC-ee*),
→ potential first step
- pp -collider (*FCC-hh*)
→ long-term goal, defining infrastructure requirements

~16 T \Rightarrow 100 TeV pp in 100 km

- HE-LHC with *FCC-hh* technology
- Ions and lepton-hadron options with hadron colliders

FCC-ee:

- **~20-50 fold improved precision on many EW quantities (equiv. to factor 5-7 in mass)** (m_Z , m_W , m_{top} , $\sin^2 \theta_W^{eff}$, R_b , $\alpha_{QED}(m_Z)$, $\alpha_s(m_Z)$, m_Z , m_W , m_τ), Higgs and top quark couplings)
- **Exploring 10 - 100 TeV energy scale via couplings with precision measurements**
- **Machine design for highest luminosities at Z, WW, ZH and ttbar working points**

FCC-hh:

- **Highest center of mass energy for direct production up to 20 - 30 TeV**
- **Huge rates for single and multiple production of SM bosons (H,W,Z) and quarks**
- **Machine design for ~100 TeV c.m. energy & int. luminosity ~ 20ab⁻¹ in 25 years**

HE-LHC:

- **Doubling LHC collision energy with FCC-hh 16 T magnet technology**
- **c.m. energy ~ 27 TeV = 14 TeV x 16 T/8.33T, target luminosity ≥ 4 x HL-LHC**
- **Machine design within constraints from LHC CE and using HL-LHC and FCC techn.**

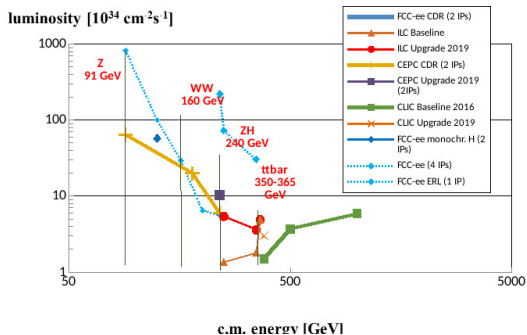
| working point | assumed typical luminosity/IP [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$] = design value minus 15(10)% | total luminosity (2 IPs)/ yr; half of typical luminosity assumed in 1st two years (Z) and 1st year ($t\bar{t}$) | physics goal | run time [yr] |
|------------------------|----------------------------------------------------------------------------------------------------------------|-------------------------------------------------------------------------------------------------------------------------------|--------------------------------------------|---------------------|
| Z first 2 years | 100 | 26 $\text{ab}^{-1}/\text{year}$ | 150 ab^{-1} | 4 |
| Z later | 200 | 48 $\text{ab}^{-1}/\text{year}$ | | |
| W | 25 | 6 $\text{ab}^{-1}/\text{year}$ | 10 ab^{-1} | 1-2 |
| H | 7.0 | 1.7 $\text{ab}^{-1}/\text{year}$ | 5 ab^{-1} | 3 |

machine modification for RF installation & rearrangement: 1 year

| | | | | |
|-------------------------------|-----|-----------------------------------------------------|----------------------------------------|----------|
| top 1st year (350 GeV) | 0.8 | 0.2 $\text{ab}^{-1}/\text{year}$ | 0.2 ab^{-1} | 1 |
| top later (365 GeV) | 1.4 | 0.34 $\text{ab}^{-1}/\text{year}$ | 1.5 ab^{-1} | 4 |

total program duration: 15 years – incl. machine modifications
phase 1 (Z, W, H): 9 years, phase 2 (top): 6 years

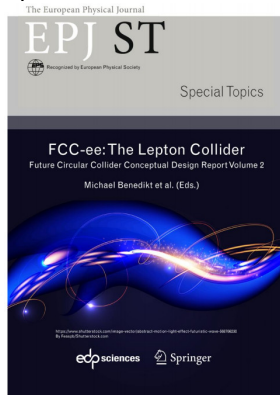
FCCee in context



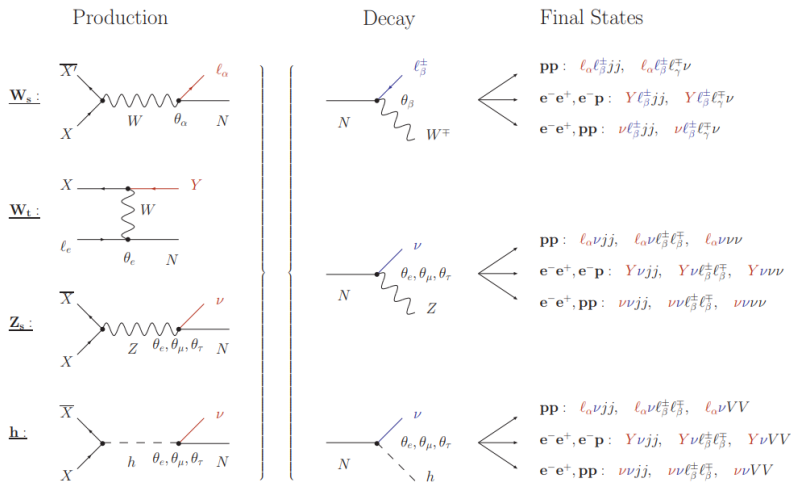
⇒ Check out the CDR : CERN-ACC-2018-0057

⇒ Also the theory report: arXiv:1905.05078

⇒ The FCCee is the most efficient machine up to the $t\bar{t}$ threshold.

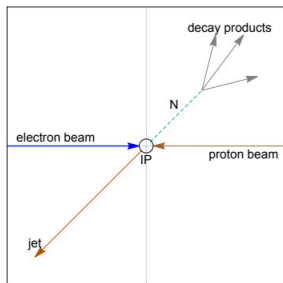
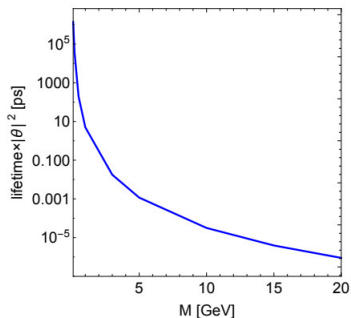


Schematizing sterile neutrino searches at FCC



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

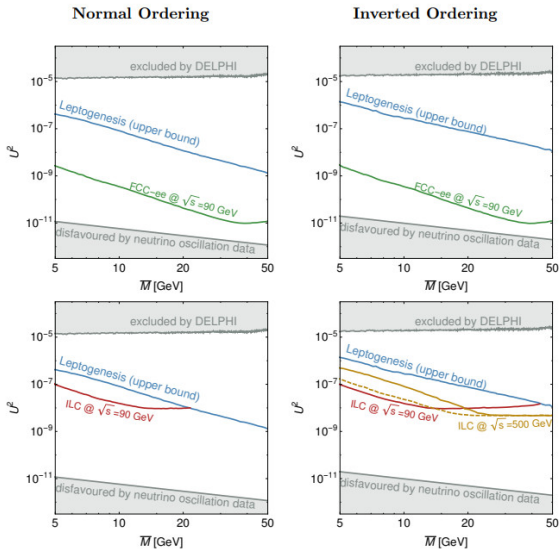
Displaced vertexes



Example: FCC-he

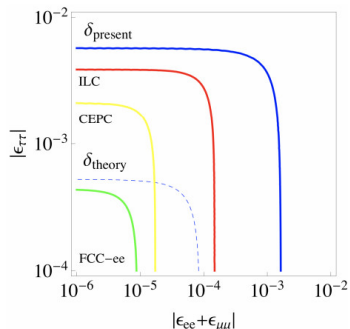
- ⇒ In the interesting region: $m < m_W$ and $\theta < 10^{-5}$
- ⇒ Displacement: measurement of primary (production) vertex.
- ⇒ Secondary vertex with „large” displacement
- ⇒ ee he: A few times tracking resolution: $\mathcal{O}(10)\mu\text{m}$,
- ⇒ hh: Beyond background, detector noise, pileup: $\mathcal{O}(10)\text{cm}$.

Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

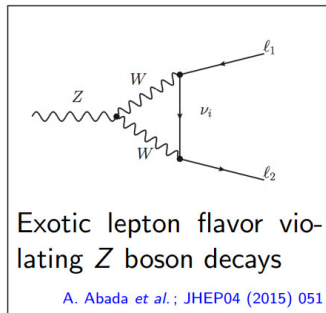


arXiv::1710.03744, S.Antuscha, E.Cazzato, M.Drewes, O.Fischer, B.Garbrecht, D.Gueterb, I.Klari

FCCee indirect

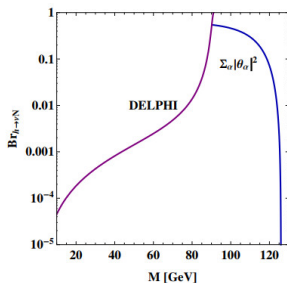
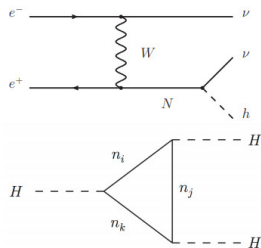


Antusch, OF; JHEP 1410 (2014) 094



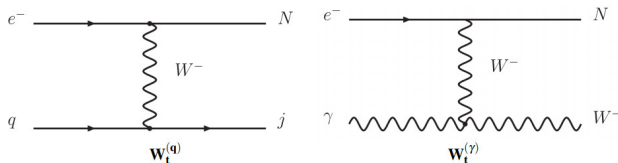
- ⇒ Modification of the theory prediction of precision observables.
- ⇒ Also CKM unitarity, cLFV, LUV.
- ⇒ Currently still dominated by LEP!

FCCee Higgs portal



S. Antusch, OF; JHEP 1604 (2016) 189

- ⇒ Mono-Higgs production mechanism!
- ⇒ New Higgs decays:
 - Modification of Higgs Branching fractions.
 - New decays: $H \rightarrow N\nu$.
 - Invisible width modification.
- ⇒ Modification of triple Higgs coupling.

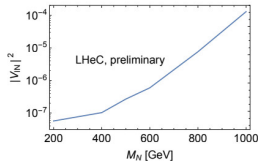
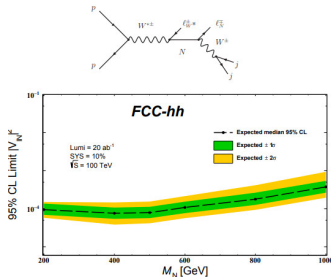


- ⇒ Large Lorentz boost makes the displaced vertexes clearly visible.
 ⇒ Many final states to look at:

| Name | Final State | $ \theta_\alpha $ Dependency | LFV |
|---------------|---------------------------------|--------------------------------------------------|-----|
| lepton-trijet | $jjj\ell_\alpha^-$ | $\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ | ✓ |
| jet-dilepton | $j\ell_\alpha^-\ell_\beta^+\nu$ | $\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*) | ✓ |
| trijet | $jjj\nu$ | $ \theta_e ^2$ | × |
| monojet | $j\nu\nu\nu$ | $ \theta_e ^2$ | × |

| Name | Final State | $ \theta_\alpha $ Dependency | LFV |
|-------------------|------------------------------------------------|--------------------------------------------------|-----|
| lepton-quadrjet | $jjjj\ell_\alpha^-$ | $\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ | ✓ |
| dilepton-dijet | $\ell_\alpha^-\ell_\beta^+\nu jj$ | $\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*) | ✓ |
| trilepton | $\ell_\alpha^-\ell_\beta^-\ell_\gamma^+\nu\nu$ | $\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*) | ✓ |
| quadrjet | $jjjj\nu$ | $ \theta_e ^2$ | × |
| electron-di-b-jet | $e^-bb\nu\nu$ | $ \theta_e ^2$ | × |
| dijet | $jj\nu\nu\nu$ | $ \theta_e ^2$ | × |
| monolepton | $\ell_\alpha^-\nu\nu\nu\nu$ | $ \theta_e ^2$ | × |

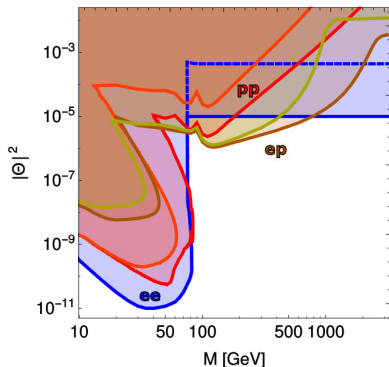
Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv.:1612.02728



S. Antusch, E. Cazzato, O. Fischer, A. Hammad and K. Wang, JHEP **1810** (2018) 067

- ⇒ LFV is the thing to look for!!!
- ⇒ The best final states: $l_{\alpha}^{\pm} l_{\beta}^{\mp} j j$, $l^{\pm} l^{\mp} l_{\gamma}^{\pm}$
- ⇒ For ep machine the more sensitive ones are: $\mu j j j$ and $\tau j j j$.
- ⇒ Also LNU are there: $\mu^{\pm} \mu^{\pm} j$ (pp) and $e^{+} j$ (ep).

FCC in total



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

⇒ FCCee:

- Dominates the exclusion below the m_W mass.
- Precision indirect constraints: EWPO, CKM, etc.

⇒ FCCeh, FCChh:

- Sensitivity in high mass region.
- Higgs potential.
- LFV, LNV.

Summary

- ⇒ Hunting for RHN is very well motivated.
- ⇒ Neutrino program has to be considered a core of future colliders.
- ⇒ FCC has unique sensitivity for RHN!
- ⇒ Huge amount of measurements and constraints to be performed.
- ⇒ Complementarity between different colliders.

Credit to M. Drewers, arXiv::1903.06100

