Extracting angular observables with the Method of Moments

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in collaboration with
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1. Motivation.
2. Method of Moments.
3. Systematic uncertainties.
4. MC toy studies.
5. Conclusions.

Motivation

Likelihood(LL) fits even though widely used suffer from couple of draw backs:

1. In case of small number events LL fits suffer from convergence problems. This behaviour is well known and was observed several times in toys for $B \to K^{*}\mu\mu$.
2. LL can exhibit a bias when underlying physics model is not well known, incomplete or mismodeled.
3. The LL have problems converging when parameters of the p.d.f. are close to their physical boundaries.
4. Accessing uncertainty in LL fits sometimes requires application of computationally expensive Feldman-Cousins method.

Method of Moments

MoM addresses the above problems:

Advantages of MoM

- Probability distribution function rapidity converges towards the Gaussian distribution.
- MoM gives an unbias result even with small data sample.
- Insensitive to large class of remodelling of physics models.
- Is completely insensitive to boundary problems.

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- "For each observable, the mean value can be determined independently from all other observables.
- Uncertainly follows perfectly $1/\sqrt{N}$ scaling, where N is number of signal events.

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Drawback:

Advantages of MoM

- Estimated uncertainty in MoM is larger then the ones from LL.

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Introduction to MoM

Let us a define a probability density function p.d.f. of a decay:

$$P(\vec{\nu},\vec{\vartheta}) \equiv \sum_i S_i(\vec{\nu}) \times f_i(\vec{\vartheta}) \quad (1)$$

Let's assume further that there exist a dual basis: $\{f_i(\vec{\vartheta})\}$, $\{\tilde{f}_i(\vec{\vartheta})\}$ that the orthogonality relation is valid:

$$\int_\Omega \mathrm{d}\vec{\vartheta}\, \tilde{f}_i(\vec{\vartheta}) f_j(\vec{\vartheta}) = \delta_{ij} \quad (2)$$

Since we want to use MoM to extract angular observables it's normal to work with Legendre polynomials. In this case we can find self-dual basis:

$$\forall_i \tilde{f}_i = f_i \ , \quad (3)$$

just by applying the ansatz: $\tilde{f}_i = \sum_i a_{ij} f_j$.

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Determination of angular observables

Thanks to the orthonormality relation Eq. 2 one can calculate the $S_i(\vec{\nu})$ just by doing the integration:

$$S_i(\vec{\nu}) = \int_{\Omega} d\vec{\vartheta} P(\vec{\nu}, \vec{\vartheta}) \tilde{f}_i(\vec{\vartheta}) \tag{4}$$

We also need to integrate out the $\vec{\nu}$ dependence:

$$\langle S_i \rangle = \int_{\Theta} d\vec{\nu} \int_{\Omega} d\vec{\vartheta} P(\vec{\nu}, \vec{\vartheta}) \tilde{f}_i(\vec{\vartheta}) \tag{5}$$

MoM is basically performing integration in Eq. 5 using MC method:

$$E[S_i] \to \widehat{E[S_i]} = \frac{1}{N} \sum_{k=1}^{N} \tilde{f}(x_k)$$

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Uncertainty estimation

MoM provides also a very fast and easy way of estimating the statistical uncertainty:

$$\sigma(S_i) = \sqrt{\frac{1}{N-1}\sum_{k=1}^{N}(\tilde{f}_i(x_k) - \widehat{S}_i)^2} \quad (6)$$

and the covariance:

$$\mathrm{Cov}[S_i, S_j] = \frac{1}{N-1}\sum_{k=1}^{N}[\widehat{S}_i - \tilde{f}_i(x_k)][\widehat{S}_j - \tilde{f}_j(x_k)] \quad (7)$$

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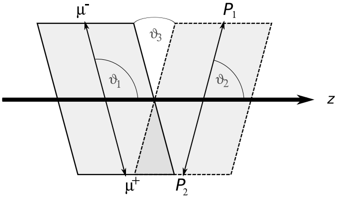
Partial Waves mismodeling

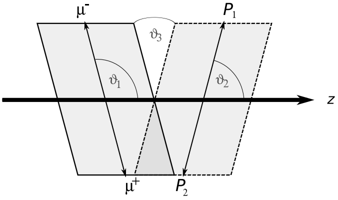
- Let us consider a decay of $B \to P_1 P_2 \mu^- \mu^+$.
- In terms of angular p.d.f. is expressed in terms of partial-wave expansion.
- For $B \to K\pi\mu^-\mu^+$ system, S,P,D waves have been studied.
  - The muon system of this kind of decays has a fixed angular dependence in terms of ϑ_1 (lepton helicity angle) and ϑ_3 (azimuthal angle).
  - The hadron system can have arbitrary large angular momentum.



Partial Waves mismodeling

- One can write the p.d.f. separating the hadronic system:

$$P(\cos\vartheta_1, \cos\vartheta_2, \vartheta_3) = \sum_i S_i(\vec{\nu}, \cos\vartheta_2) f_i(\cos\vartheta_1, \vartheta_3) \quad (8)$$



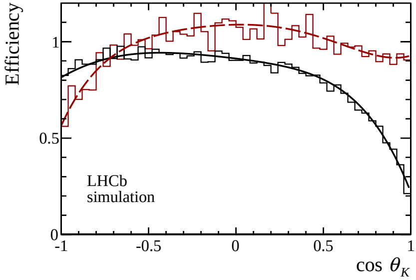
  - $S_i(\vec{\nu}, \cos\vartheta_2)$ can be further expend in terms of Legendre polynomials $p_l^{|m|}(\cos\vartheta_2)$:

$$S_i(\vec{\nu}, \cos\vartheta_2) = \sum_{l=0}^{\inf} S_{k,l}(\vec{\nu}) p_l^{|m|}(\cos\vartheta_2) \quad (9)$$

  - Experimentally the $S_{k,l}$ are easily accessible, but there is a theoretical difficulty as one would need to sum over infinite number of partial waves.

Detector effects

- Since our detectors are not a perfect devices the angular distribution observed by them are not the distributions that the physics model creates.



- To take into account the acceptance effects one needs to simulate the a large sample of MC events.
- Try to figure out the efficiency function.
- Number of possibilities.
- Then you can just weight events:

$$\widehat{E[S_i]} = \frac{1}{\sum_{k=1}^{N} w_k} \sum_{k=1}^{N} w_k \tilde{f}(x_k), \;\; w_k = \frac{1}{\epsilon(x_k)}$$

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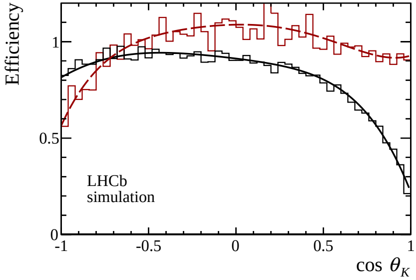

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Unfolding matrix

In general one can write the distribution of events after the detector effects:

$$P^{\mathrm{Det}}(x_d) = N \int \int dx_t \; P^{\mathrm{Phys}}(x_t) E(x_d|x_t), \qquad (10)$$

where $N^{-1} = \int \int dx_t \; dx_d \; P^{\mathrm{Phys}}(x_t) E(x_d|x_t)$ and $E(x_d|x_t)$ denotes the efficiency $\epsilon(x_t)$ and resolution of the detector $R(x_d|x_t)$:

$$E(x_d|x_t) = \epsilon(x_t) R(x_d|x_t) \qquad (11)$$

One can define the raw moments:

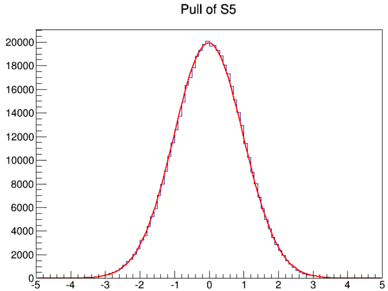
$$Q_i^{(m)} = \int \int dx_t \; dx_d \; \tilde{f}_i(x_d) P^{(m)}(x_t) E(x_d|x_t) \qquad (12)$$

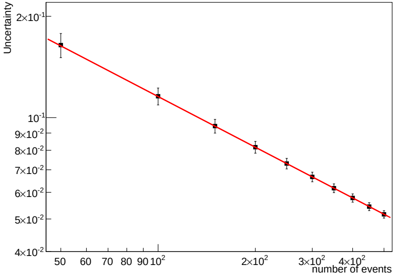
$$M_{ij} = \begin{cases} 2 Q_i^{(0)} & j = 0\,, \\ 2\left(Q_i^{(j)} - Q_i^{(0)}\right) & j \neq 0\,, \end{cases} \qquad (13)$$

Once we measured the moments Q in data we can invert Eq. 11 and get the $\vec{S}$: $\widehat{\vec{S}} = M^{-1} \widehat{\vec{Q}}$.

Toy Validation

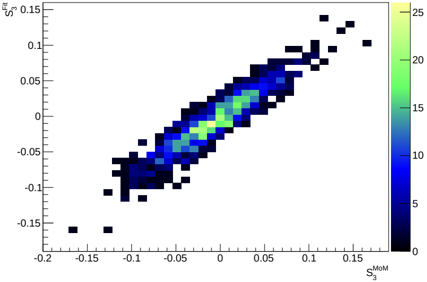
- All the statistics properties of MoM have been tested in numbers of TOY MC.
- As long as you have ~ 30 events your pulls are perfectly gaussian.
- Uncertainty scales with $\frac{\alpha}{\sqrt{n}}$, $\alpha = \mathcal{O}(1)$.
- Never observed any boundary problems.

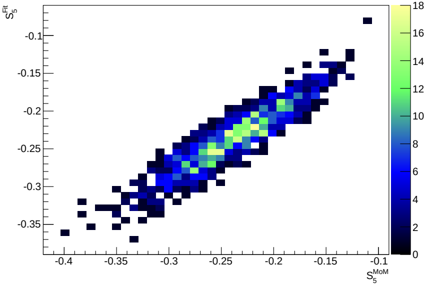


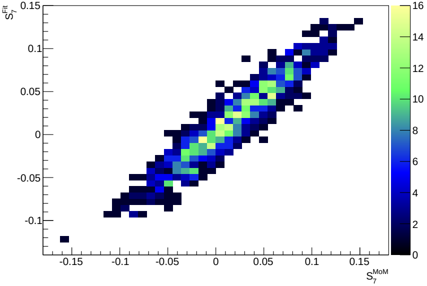


Correlation of MoM with Likelihood

- MoM is highly correlated with LL.
- Despite the correlation there can be difference of the order of statistical error.







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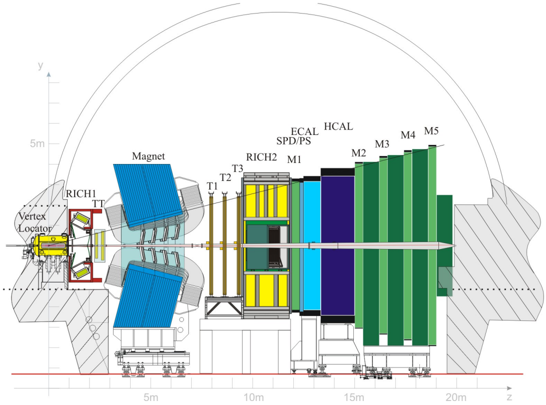
Conclusions

1. MoM viable alternative to LL fits.
2. Allows LHCb to go smaller q^2 bins (get ready for 1 GeV^2 soon!).
3. Alternative method of extracting the detector effects.
4. Method is universally applicable, as long as an orthonormal basis for the p.d.f. exists.

BACKUP

LHCb detector



LHCb is a forward spectrometer:

- Excellent vertex resolution.
- Efficient trigger.
- High acceptance for τ and B.
- Great Particle ID

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