

# Searches for heavy neutral leptons at the Future Circular Colliders



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on behalf of the FCC collaboration,  
CERN and

Institute of Nuclear Physics, Polish Academy of Science

15th Rencontres du Vietnam,  
Quy Nhon, 4-10 August 2019

# Particles of SM

## SM

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	<b>u</b> Left up Right	<b>c</b> Left charm Right	<b>t</b> Left top Right
Quarks	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	<b>d</b> Left down Right	<b>s</b> Left strange Right	<b>b</b> Left bottom Right
Leptons	0 eV	0 eV	0 eV
	0	0	0
	<b><math>\nu_e</math></b> Left electron neutrino Right	<b><math>\nu_\mu</math></b> Left muon neutrino Right	<b><math>\nu_\tau</math></b> Left tau neutrino Right
0.511 MeV	105.7 MeV	1.777 GeV	
-1	-1	-1	
<b>e</b> Left electron Right	<b><math>\mu</math></b> Left muon Right	<b><math>\tau</math></b> Left tau Right	

# Particles of SM

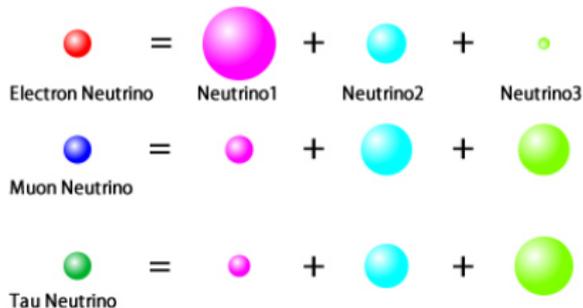
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	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
<b>Quarks</b>	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom
	Left Right	Left Right	Left Right
	0 eV	0 eV	0 eV
	$0$	$0$	$0$
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino
	Left Right	Left Right	Left Right
	0.511 MeV	105.7 MeV	1.777 GeV
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<b>Leptons</b>	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau
	Left Right	Left Right	Left Right

⇒ In the SM neutrinos do not appear in the right-handed state.

⇒ By construction neutrinos are massless.

⇒ Neutrino oscillations are evidence for physics beyond the SM!



# Extending the SM

⇒ A lazy person solution is to add the right-handed neutrinos in:

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	Left Right	Left Right	Left Right
	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
Quarks	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom
	Left Right	Left Right	Left Right
	0 eV	0 eV	0 eV
	$0$	$0$	$0$
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino
	Left Right	Left Right	Left Right
	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
Leptons	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau
	Left Right	Left Right	Left Right

**nuMSM**

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	<b>u</b> up	<b>c</b> charm	<b>t</b> top
	Left Right	Left Right	Left Right
	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
Quarks	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom
	Left Right	Left Right	Left Right
	<0.0001 eV	~10 keV	~GeV
	$0$	$0$	$0$
	<b><math>\nu_e</math></b> <b><math>N_1</math></b> electron neutrino sterile neutrino	<b><math>\nu_\mu</math></b> <b><math>N_2</math></b> muon neutrino sterile neutrino	<b><math>\nu_\tau</math></b> <b><math>N_3</math></b> tau neutrino sterile neutrino
	Left Right	Left Right	Left Right
	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
Leptons	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau
	Left Right	Left Right	Left Right

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SM						nuMSM									
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charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	name →	Left	Right	Left	Right	Left	Right	Left
name →	Left	Right	Left	name →	Left	Right	Left	name →	Left	Right	Left	Right	Left	Right	Left
	<b>u</b>	<b>c</b>	<b>t</b>		<b>u</b>	<b>c</b>	<b>t</b>		up	charm	top		up	charm	top
	4.8 MeV	104 MeV	4.2 GeV		4.8 MeV	104 MeV	4.2 GeV		down	strange	bottom		down	strange	bottom
Quarks	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	Quarks	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$		down	strange	bottom		down	strange	bottom
	<b>d</b>	<b>s</b>	<b>b</b>		<b>d</b>	<b>s</b>	<b>b</b>		0 eV	0 eV	0 eV		$<0.0001$ eV	$\sim 10$ keV	$\sim 0.01$ eV
	0 eV	0 eV	0 eV		$<0.0001$ eV	$\sim 10$ keV	$\sim 0.01$ eV		$0 \nu_e$	$0 \nu_\mu$	$0 \nu_\tau$		$0 \nu_e N_1$	$0 \nu_\mu N_2$	$0 \nu_\tau N_3$
	electron neutrino	muon neutrino	tau neutrino		electron neutrino	muon neutrino	tau neutrino		electron neutrino	muon neutrino	tau neutrino		electron neutrino	muon neutrino	tau neutrino
	0.511 MeV	105.7 MeV	1.777 GeV		0.511 MeV	105.7 MeV	1.777 GeV		sterile neutrino	sterile neutrino	sterile neutrino		sterile neutrino	sterile neutrino	sterile neutrino
Leptons	-1	-1	-1	Leptons	-1	-1	-1		electron	muon	tau		electron	muon	tau
	<b>e</b>	<b><math>\mu</math></b>	<b><math>\tau</math></b>		<b>e</b>	<b><math>\mu</math></b>	<b><math>\tau</math></b>		electron	muon	tau		electron	muon	tau

⇒ But where are they?

arXiv::hep-ph/0605047, M.Shaposhnikov

# Seesaw mechanism

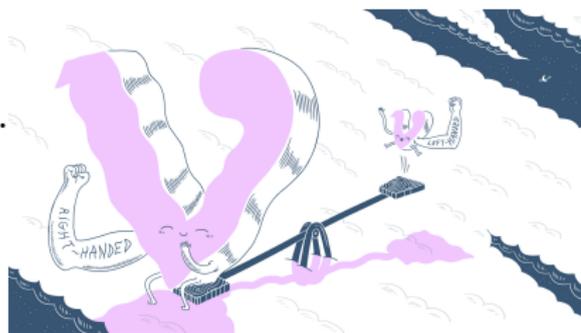
$$\mathcal{L} = \mathcal{L}_{SM} + \bar{\ell}_L F \nu_R \epsilon \Phi^* - \frac{1}{2} \bar{\nu}_R^c M_M \nu_R + \text{H.c.}$$

⇒ After the EWSB:

$$\frac{1}{2} (\bar{\nu}_L \bar{\nu}_R^c) \mathcal{M} (\nu_L^c \nu_R)^T$$

⇒ In the vanilla seesaw:

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_M \end{pmatrix} \quad \Rightarrow \lambda_+ \sim M_D, \quad \lambda_- \sim -\frac{M_M^2}{M_D}$$



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⇒ In reality:

$$\mathcal{M} = \begin{pmatrix} \delta m_\nu^{1loop} & M_D \\ M_D^T & M_M + \delta M_N^{1loop} \end{pmatrix}$$



# Seesaw mechanism

⇒ Diagonalization matrix:

$$U = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta^\dagger) & \cos(\theta^\dagger) \end{pmatrix} \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix}$$

$$U^\dagger M U^* = \begin{pmatrix} m_\nu^{\text{diag}} & \\ & M_N^{\text{diag}} \end{pmatrix}$$

with

$$M_N^{\text{diag}} = U_N^T M_N U_N = \text{diag}(M_1, M_2, M_3)$$

$$m_\nu^{\text{diag}} = U_\nu^\dagger m_\nu U_\nu^* = \text{diag}(m_1, m_2, m_3).$$

⇒ For small mixings:

$$U = \left[ \begin{pmatrix} \mathbb{I} - \frac{1}{2}\theta\theta^\dagger & \theta \\ -\theta^\dagger & \mathbb{I} - \frac{1}{2}\theta^\dagger\theta \end{pmatrix} + \mathcal{O}(\theta^3) \right] \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix},$$

# Correction to SM processes

⇒ Charge currents:

$$j_{\mu}^{+} = \frac{g}{2} \theta_{\alpha} \bar{\ell}_{\alpha} \gamma_{\mu} N$$

⇒ Neutral currents:

$$j_{\mu}^0 = \nu_{\alpha} \gamma_{\mu} \theta_{\alpha} N$$

⇒ The Yukawa couplings:

$$\mathcal{L}_{Yukawa} = \mathcal{L}_{Yukawa}^{SM} \theta_{\alpha}$$

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

## Indirect

- EW precision observables
- LFV, LNV
- Neutrinoless double beta decay
- Big Bang Nucleosynthesis

# Current status

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

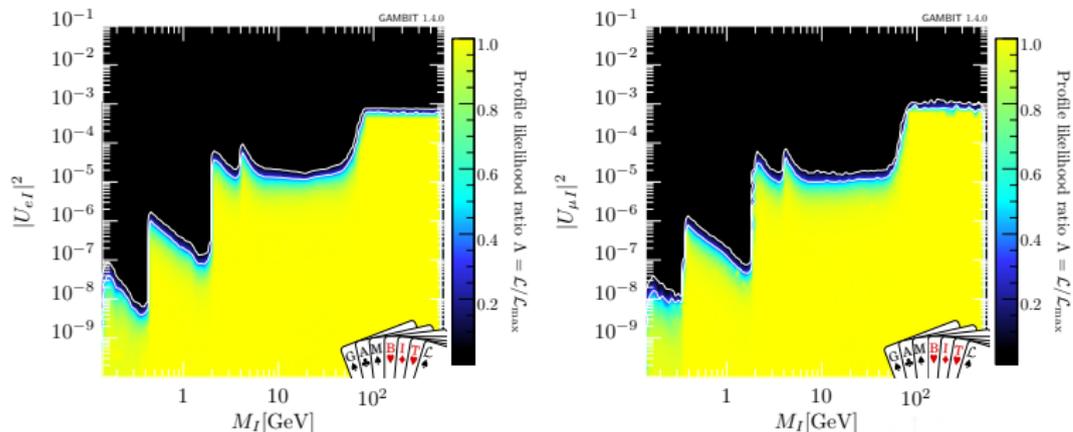
## Indirect

- EW precision observables
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- Neutrinoless double beta decay
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## Direct

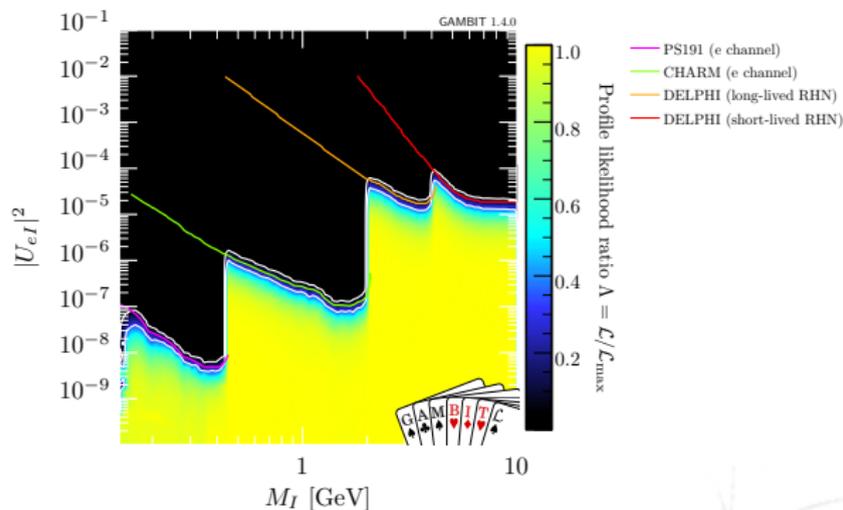
- Fix target experiments
- Collider searches (LEP, LHC, etc.)

# Current status



⇒ M.C., M. Drewers, T. Gonzalo, J. Harz, S. Krishnamurthy, C. Weniger,  
arXiv::1908.02302

# Where the constrains come from?

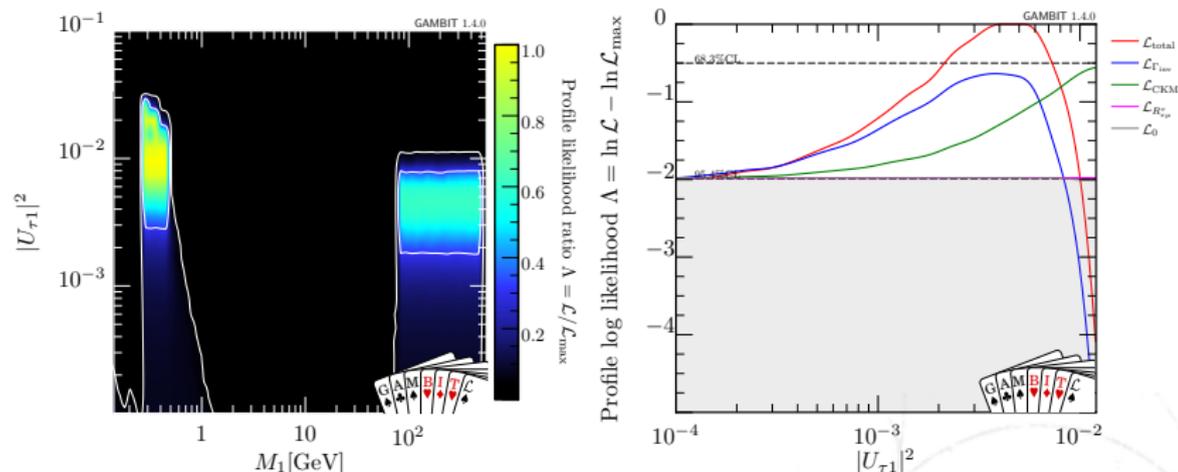


⇒ The direct searches are the strongest constraints where production cross sections are the largest.

⇒ arXiv::1908.02302

# Indirect constraints

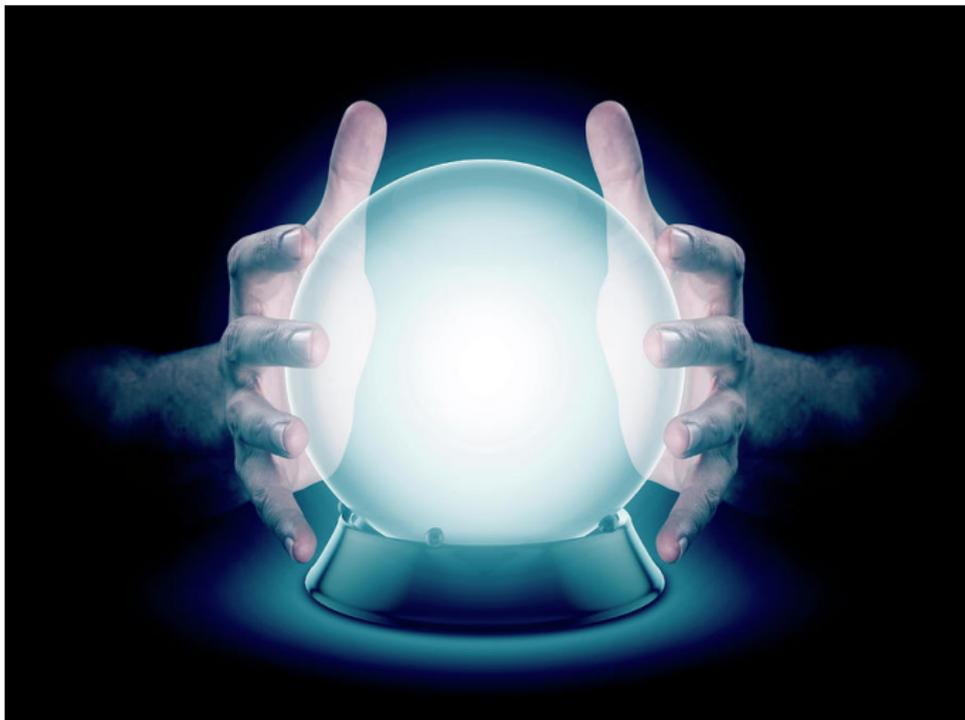
⇒ The indirect searches show power for high couplings:



⇒ Small excess is visible. Consistent with the fluctuation.

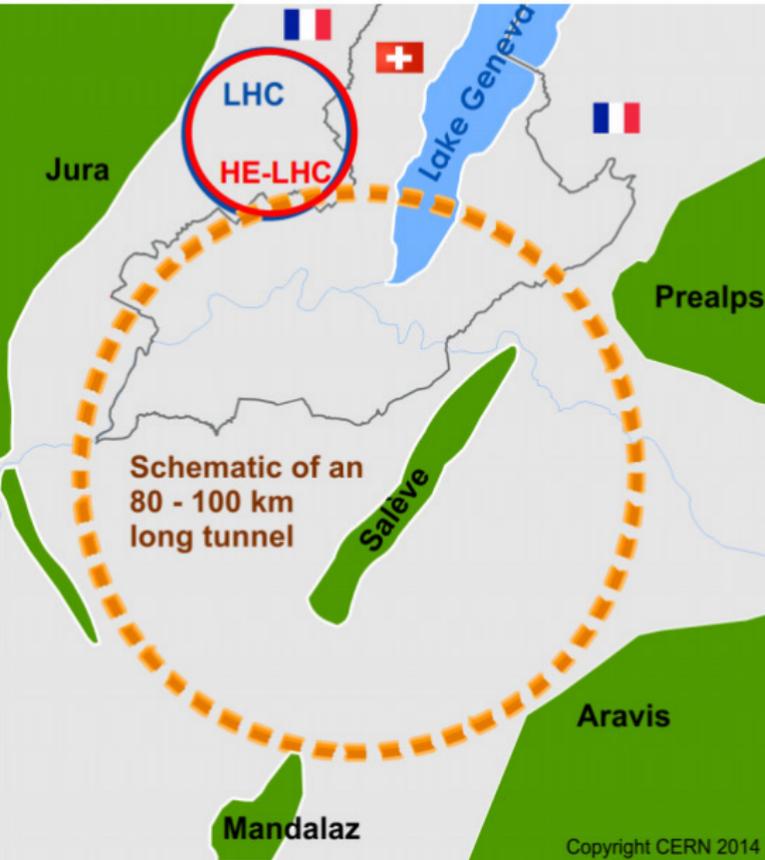
⇒ arXiv::1908.02302

What can happen in the future?



What can happen in the future?





## International FCC collaboration with CERN as host lab to study:

- ~100 km tunnel infrastructure in Geneva area, linked to CERN
- $e^+e^-$  collider (*FCC-ee*),  
→ potential first step
- $pp$ -collider (*FCC-hh*)  
→ long-term goal, defining infrastructure requirements

**~16 T  $\Rightarrow$  100 TeV  $pp$  in 100 km**

- **HE-LHC** with *FCC-hh* technology
- **ions** and **lepton-hadron** options with hadron colliders

## FCC-ee:

- **~20-50 fold improved precision on many EW quantities (equiv. to factor 5-7 in mass)** ( $m_Z$ ,  $m_W$ ,  $m_{top}$ ,  $\sin^2 \theta_W^{eff}$ ,  $R_b$ ,  $\alpha_{QED}(m_Z)$ ,  $\alpha_s(m_Z)$ ,  $m_Z$ ,  $m_W$ ,  $m_\tau$ ), Higgs and top quark couplings)
- **Exploring 10 - 100 TeV energy scale via couplings with precision measurements**
- **Machine design for highest luminosities at Z, WW, ZH and  $t\bar{t}$  working points**

## FCC-hh:

- **Highest center of mass energy for direct production up to 20 - 30 TeV**
- **Huge rates for single and multiple production of SM bosons (H,W,Z) and quarks**
- **Machine design for ~100 TeV c.m. energy & int. luminosity ~ 20ab<sup>-1</sup> in 25 years**

## HE-LHC:

- **Doubling LHC collision energy with FCC-hh 16 T magnet technology**
- **c.m. energy ~ 27 TeV = 14 TeV x 16 T/8.33T, target luminosity  $\geq 4$  x HL-LHC**
- **Machine design within constraints from LHC CE and using HL-LHC and FCC techn.**

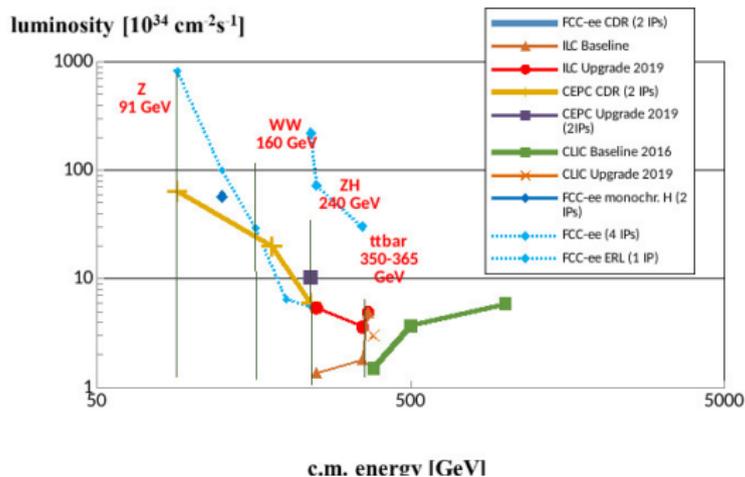
working point	assumed typical luminosity/IP [ $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ ] = design value minus 15(10)%	total luminosity (2 IPs)/ yr; half of typical luminosity assumed in 1st two years (Z) and 1st year ( $t\bar{t}$ )	physics goal	run time [yr]
<b>Z first 2 years</b>	100	<b>26 <math>\text{ab}^{-1}/\text{year}</math></b>	<b>150 <math>\text{ab}^{-1}</math></b>	<b>4</b>
<b>Z later</b>	200	<b>48 <math>\text{ab}^{-1}/\text{year}</math></b>		
<b>W</b>	25	<b>6 <math>\text{ab}^{-1}/\text{year}</math></b>	<b>10 <math>\text{ab}^{-1}</math></b>	<b>1-2</b>
<b>H</b>	7.0	<b>1.7 <math>\text{ab}^{-1}/\text{year}</math></b>	<b>5 <math>\text{ab}^{-1}</math></b>	<b>3</b>

machine modification for RF installation & rearrangement: 1 year

<b>top 1st year (350 GeV)</b>	0.8	<b>0.2 <math>\text{ab}^{-1}/\text{year}</math></b>	<b>0.2 <math>\text{ab}^{-1}</math></b>	<b>1</b>
<b>top later (365 GeV)</b>	1.4	<b>0.34 <math>\text{ab}^{-1}/\text{year}</math></b>	<b>1.5 <math>\text{ab}^{-1}</math></b>	<b>4</b>

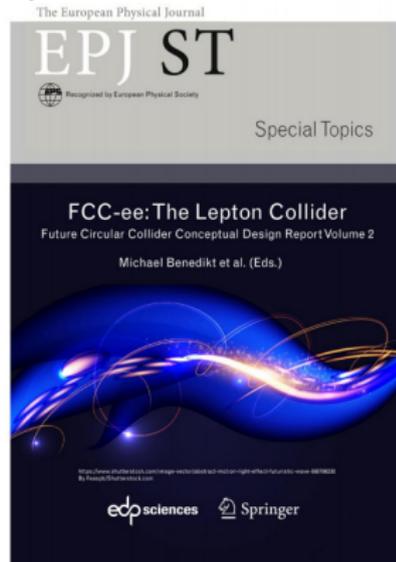
**total program duration: 15 years – incl. machine modifications**  
**phase 1 (Z, W, H): 9 years, phase 2 (top): 6 years**

# FCCee in context

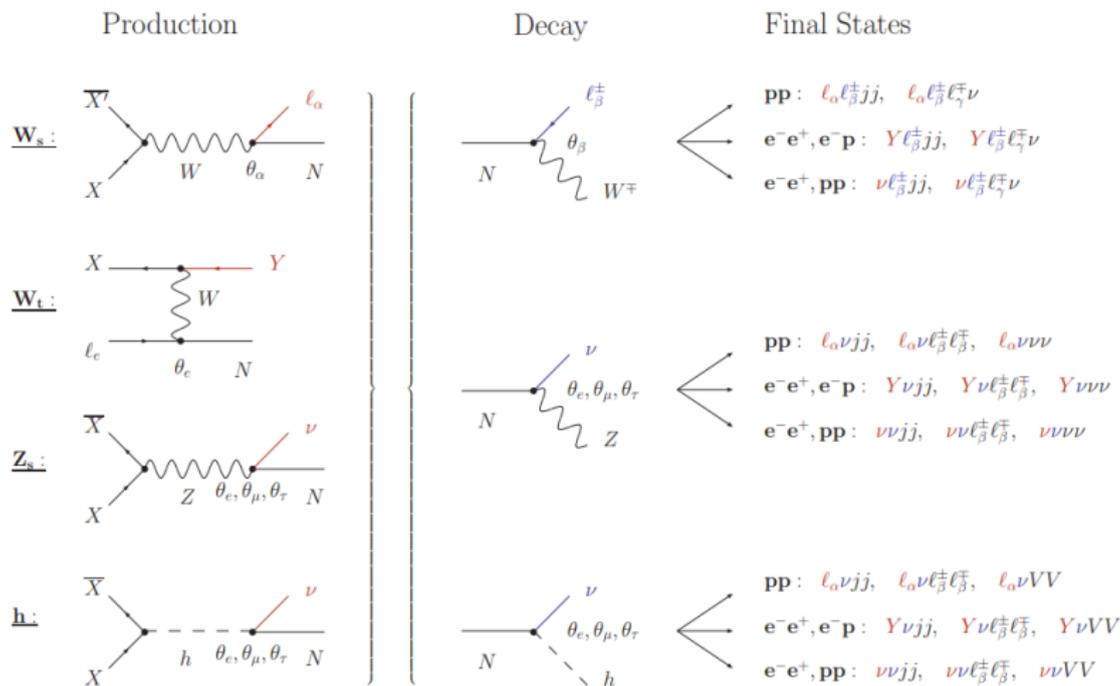


- ⇒ Check out the CDR : CERN-ACC-2018-0057
- ⇒ Also the theory report: arXiv:1905.05078

⇒ The FCCee is the most efficient machine up to the  $t\bar{t}$  threshold.

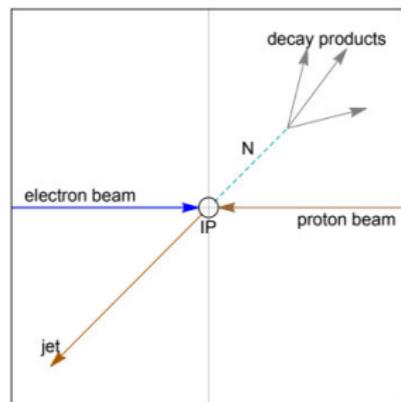
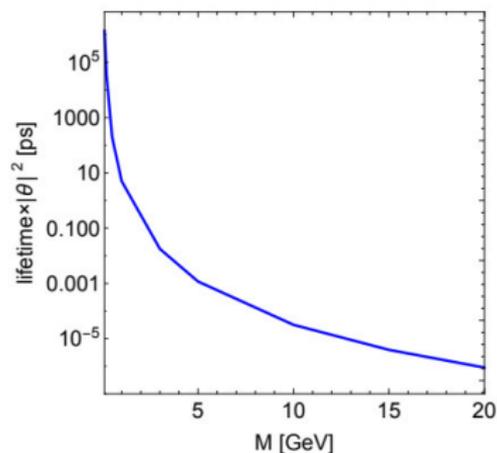


# Schematizing sterile neutrino searches at FCC



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

# Displaced vertexes

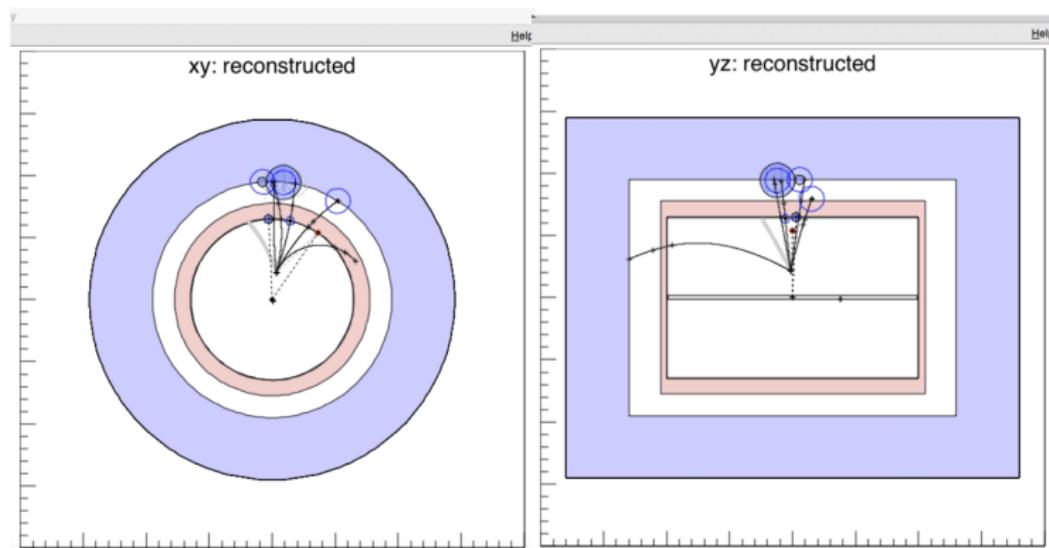


Example: FCC-he

- ⇒ In the interesting region:  $m < m_W$  and  $\theta < 10^{-5}$
- ⇒ Displacement: measurement of primary (production) vertex.
- ⇒ Secondary vertex with „large” displacement
- ⇒ ee he: A few times tracking resolution:  $\mathcal{O}(10)\mu\text{m}$ ,
- ⇒ hh: Beyond background, detector noise, pileup:  $\mathcal{O}(10)\text{cm}$ .

Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

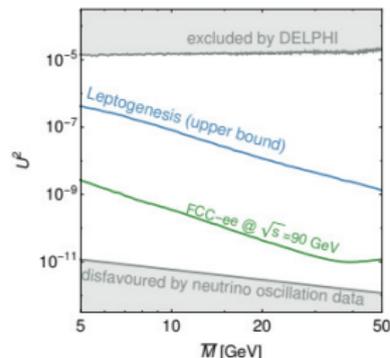
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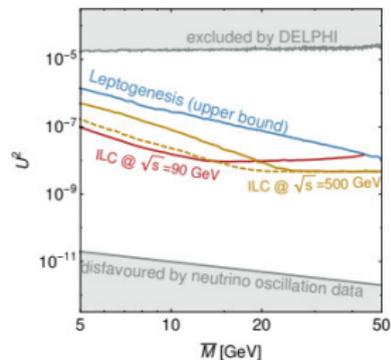
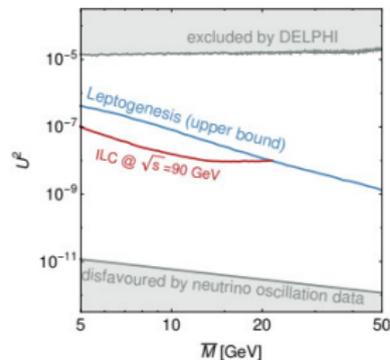
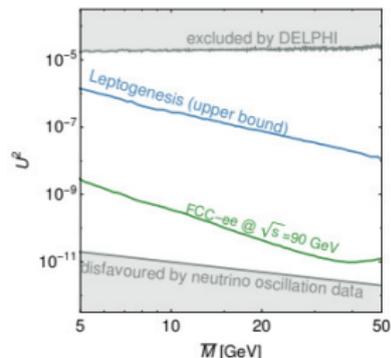
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Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv:1612.02722

## Normal Ordering

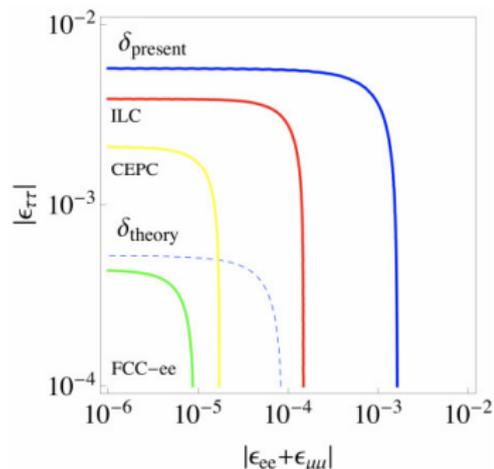


## Inverted Ordering

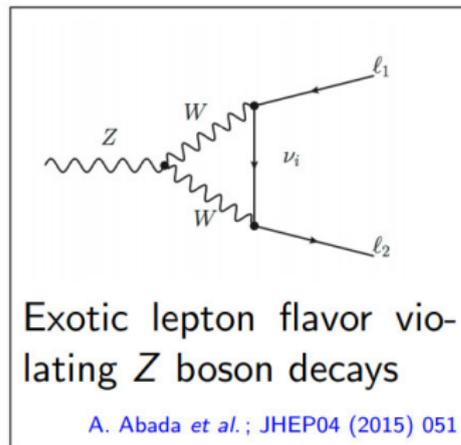


arXiv:1710.03744, S.Antusch, E.Cazzato, M.Drewes, O.Fischer, B.Garbrecht, D.Gueter, I.Klaric

# FCCee indirect

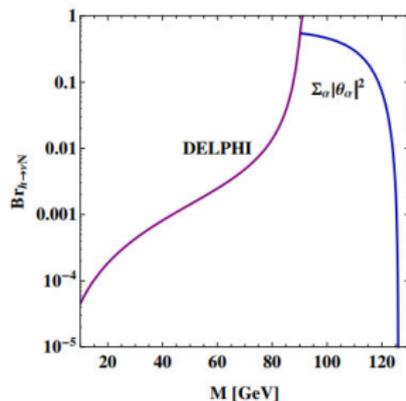
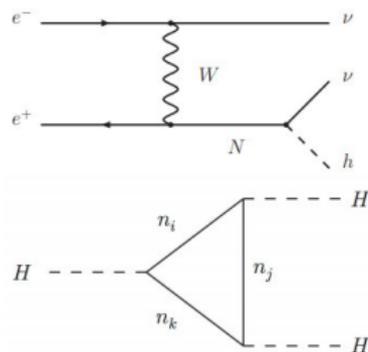


Antusch, OF; JHEP 1410 (2014) 094



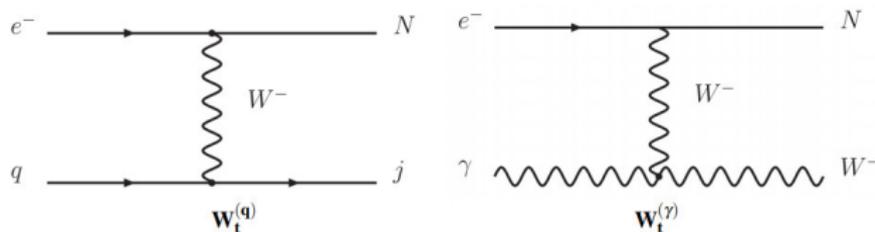
- ⇒ Modification of the theory prediction of precision observables.
- ⇒ Also CKM unitarity, cLFV, LUV.
- ⇒ Currently still dominated by LEP!

# FCCee Higgs portal



S. Antusch, OF; JHEP 1604 (2016) 189

- ⇒ Mono-Higgs production mechanism!
- ⇒ New Higgs decays:
  - Modification of Higgs Branching fractions.
  - New decays:  $H \rightarrow N\nu$ .
  - Invisible width modification.
- ⇒ Modification of triple Higgs coupling.

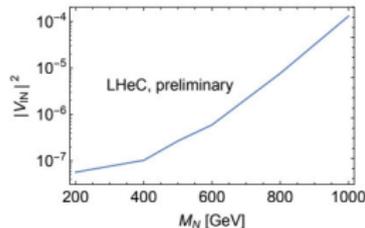
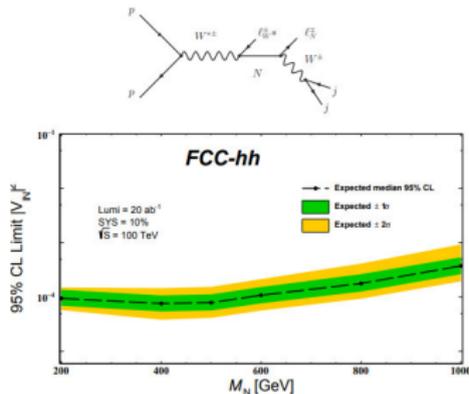


- ⇒ Large Lorentz boost makes the displaced vertexes clearly visible.  
 ⇒ Many final states to look at:

Name	Final State	$ \theta_\alpha $ Dependency	LFV
lepton-trijet	$jjj\ell_\alpha^-$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	✓
jet-dilepton	$j\ell_\alpha^-\ell_\beta^+\nu$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*)	✓
trijet	$jjj\nu$	$ \theta_e ^2$	×
monojet	$j\nu\nu\nu$	$ \theta_e ^2$	×

Name	Final State	$ \theta_\alpha $ Dependency	LFV
lepton-quadrjet	$jjjj\ell_\alpha^-$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	✓
dilepton-dijet	$\ell_\alpha^-\ell_\beta^+\nu jj$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*)	✓
trilepton	$\ell_\alpha^-\ell_\beta^-\ell_\gamma^+\nu\nu$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$ (*)	✓
quadrjet	$jjjj\nu$	$ \theta_e ^2$	×
electron-di-b-jet	$e^-bb\nu\nu$	$ \theta_e ^2$	×
dijet	$jj\nu\nu\nu$	$ \theta_e ^2$	×
monolepton	$\ell_\alpha^-\nu\nu\nu\nu$	$ \theta_e ^2$	×

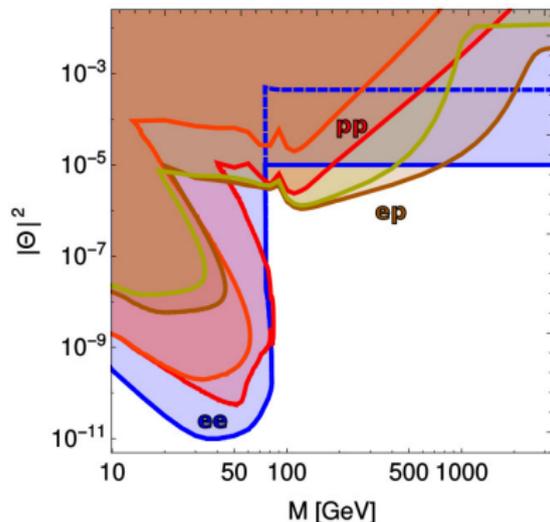
Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv.:1612.02728



S. Antusch, E. Cazzato, O. Fischer, A. Hammad and K. Wang, *JHEP* **1810** (2018) 067

- ⇒ LFV is the thing to look for!!!
- ⇒ The best final states:  $l_\alpha^\pm l_\beta^\mp jj$ ,  $l^\pm l^\mp l_\gamma^\pm$
- ⇒ For ep machine the more sensitive ones are:  $\mu jjj$  and  $\tau jjj$ .
- ⇒ Also LNU are there:  $\mu^\pm \mu^\pm j$  (pp) and  $e^+ j$  (ep).

# FCC in total



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

⇒ FCCee:

- Dominates the exclusion below the  $m_W$  mass.
- Precision indirect constraints: EWPO, CKM, etc.

⇒ FCCeh, FCChh:

- Sensitivity in high mass region.
- Higgs potential.
- LFV, LNV.

# Summary

- ⇒ Hunting for RHN is very well motivated.
- ⇒ Neutrino program has to be considered a core of future colliders.
- ⇒ FCC has unique sensitivity for RHN!
- ⇒ Huge amount of measurements and constraints to be performed.
- ⇒ Complementarity between different colliders.

Credit to M. Drewers, [Slides]

