

Searches for heavy neutral leptons at the Future Circular Colliders



Marcin Chrzaszcz
mchrzasz@cern.ch



on behalf of the FCC collaboration,
CERN and

Institute of Nuclear Physics, Polish Academy of Science

15th Rencontres du Vietnam,
Quy Nhon, 4-10 August 2019

Particles of SM

SM

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	u up	c charm	t top
	Left Right	Left Right	Left Right
Quarks	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	d down	s strange	b bottom
	Left Right	Left Right	Left Right
	0 eV	0 eV	0 eV
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino
	Left Right	Left Right	Left Right
Leptons	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
	e electron	μ muon	τ tau
	Left Right	Left Right	Left Right

Particles of SM

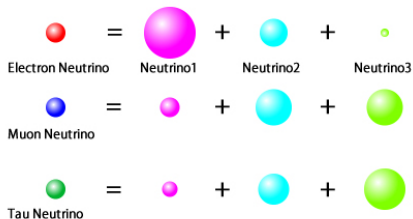
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⇒ In the SM neutrinos do not appear in the right-handed state.

⇒ By construction neutrinos are massless.

⇒ Neutrino oscillations are evidence for physics beyond the SM!



Extending the SM

⇒ A lazy person solution is to add the right-handed neutrinos in:

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Leptons	e electron	μ muon	τ tau
	Left Right	Left Right	Left Right

nuMSM

mass →	2.4 MeV	1.27 GeV	171.2 GeV
charge →	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
name →	u up	c charm	t top
	Left Right	Left Right	Left Right
	4.8 MeV	104 MeV	4.2 GeV
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
Quarks	d down	s strange	b bottom
	Left Right	Left Right	Left Right
	<0.0001 eV	~10 keV	~GeV
	0	0	0
	ν_e N_1 electron neutrino sterile neutrino	ν_μ N_2 muon neutrino sterile neutrino	ν_τ N_3 tau neutrino sterile neutrino
	Left Right	Left Right	Left Right
	0.511 MeV	105.7 MeV	1.777 GeV
	-1	-1	-1
Leptons	e electron	μ muon	τ tau
	Left Right	Left Right	Left Right

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name →	u up	c charm	t top	u up	c charm	t top	
	Left Right	Left Right	Left Right	Left Right	Left Right	Left Right	
	4.8 MeV	104 MeV	4.2 GeV	4.8 MeV	104 MeV	4.2 GeV	
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$	
Quarks	d down	s strange	b bottom	d down	s strange	b bottom	
	Left Right	Left Right	Left Right	Left Right	Left Right	Left Right	
	0 eV	0 eV	0 eV	<0.0001 eV	~ 10 keV	~ 0.01 eV	\sim GeV
	0	0	0	0	0	0	0
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	N_1 sterile neutrino
	Left Right	Left Right	Left Right	Left Right	Left Right	Left Right	Left Right
	0.511 MeV	105.7 MeV	1.777 GeV	0.511 MeV	105.7 MeV	1.777 GeV	
	-1	-1	-1	-1	-1	-1	
Leptons	e electron	μ muon	τ tau	e electron	μ muon	τ tau	
	Left Right	Left Right	Left Right	Left Right	Left Right	Left Right	

⇒ But where are they?

arXiv::hep-ph/0605047, M.Shaposhnikov

Seesaw mechanism

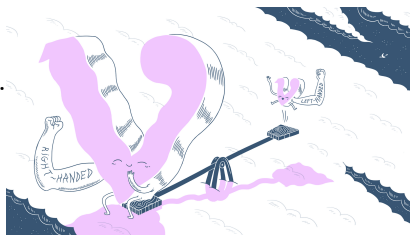
$$\mathcal{L} = \mathcal{L}_{SM} + \bar{\ell}_L F \nu_R \epsilon \Phi^* - \frac{1}{2} \bar{\nu}_R^c M_M \nu_R + \text{H.c.}$$

⇒ After the EWSB:

$$\frac{1}{2} \left(\bar{\nu}_L \bar{\nu}_R^c \right) \mathcal{M} \left(\nu_L^c \nu_R \right)^T$$

⇒ In the vanilla seesaw:

$$\mathcal{M} = \begin{pmatrix} 0 & M_D \\ M_D & M_M \end{pmatrix} \quad \Rightarrow \lambda_+ \sim M_D, \quad \lambda_- \sim -\frac{M_M^2}{M_D}$$



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⇒ In reality:

$$\mathcal{M} = \begin{pmatrix} \delta m_\nu^{1loop} & M_D \\ M_D^T & M_M + \delta M_N^{1loop} \end{pmatrix}$$



Seesaw mechanism

⇒ Diagonalization matrix:

$$U = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta^\dagger) & \cos(\theta^\dagger) \end{pmatrix} \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix}$$

$$U^\dagger M U^* = \begin{pmatrix} m_\nu^{\text{diag}} & \\ & M_N^{\text{diag}} \end{pmatrix}$$

with

$$M_N^{\text{diag}} = U_N^T M_N U_N = \text{diag}(M_1, M_2, M_3)$$

$$m_\nu^{\text{diag}} = U_\nu^\dagger m_\nu U_\nu^* = \text{diag}(m_1, m_2, m_3).$$

⇒ For small mixings:

$$U = \left[\begin{pmatrix} \mathbb{I} - \frac{1}{2}\theta\theta^\dagger & \theta \\ -\theta^\dagger & \mathbb{I} - \frac{1}{2}\theta^\dagger\theta \end{pmatrix} + \mathcal{O}(\theta^3) \right] \begin{pmatrix} U_\nu \\ U_N^* \end{pmatrix},$$

Correction to SM processes

⇒ Charge currents:

$$j_{\mu}^{+} = \frac{g}{2} \theta_{\alpha} \bar{\ell}_{\alpha} \gamma_{\mu} N$$

⇒ Neutral currents:

$$j_{\mu}^0 = \nu_{\alpha} \gamma_{\mu} \theta_{\alpha} N$$

⇒ The Yukawa couplings:

$$\mathcal{L}_{Yukawa} = \mathcal{L}_{Yukawa}^{SM} \theta_{\alpha}$$

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

Indirect

- EW precision observables
- LFV, LNV
- Neutrinoless double beta decay
- Big Bang Nucleosynthesis

Current status

⇒ Since the RHN are modifying fundamental properties of SM they are hugely constrained:

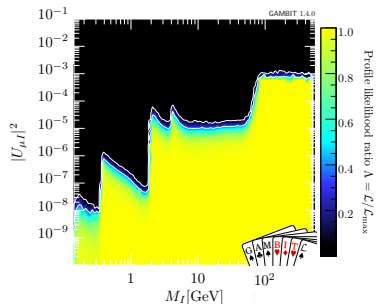
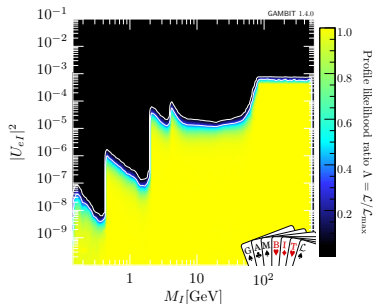
Indirect

- EW precision observables
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Direct

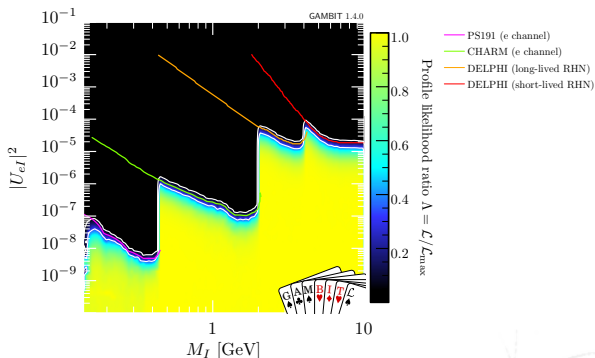
- Fix target experiments
- Collider searches (LEP, LHC, etc.)

Current status



⇒ M.C., M. Drewers, T. Gonzalo, J. Harz, S. Krishnamurthy, C. Weniger,
arXiv::1908.02302

Where the constrains come from?

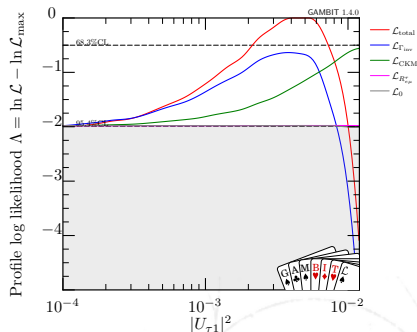
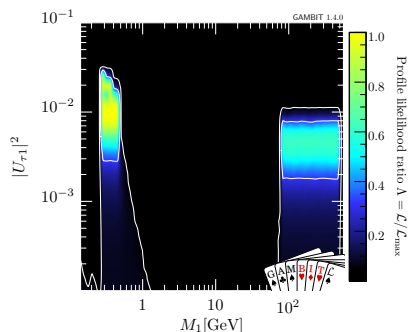


⇒ The direct searches are the strongest constraints where production cross sections are the largest.

⇒ arXiv::1908.02302

Indirect constraints

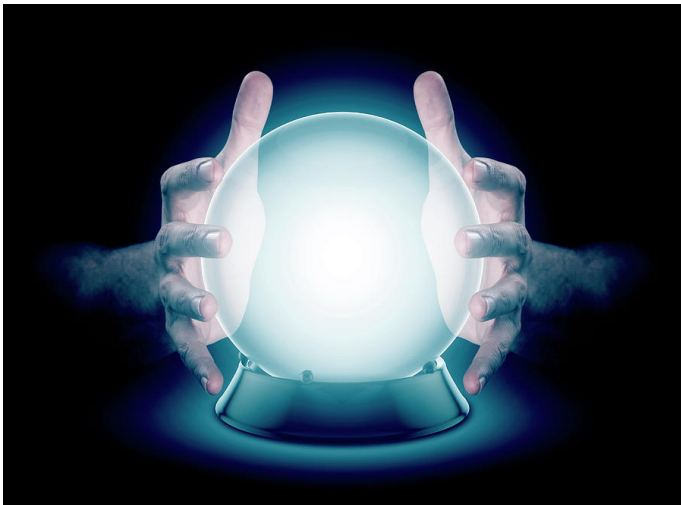
⇒ The indirect searches show power for high couplings:



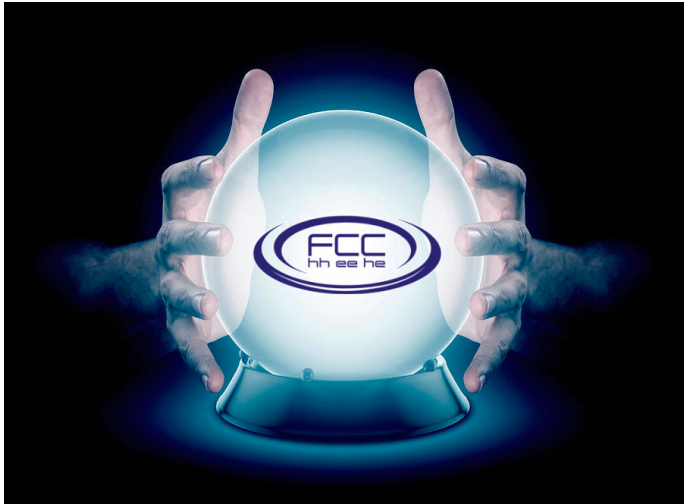
⇒ Small excess is visible. Consistent with the fluctuation.

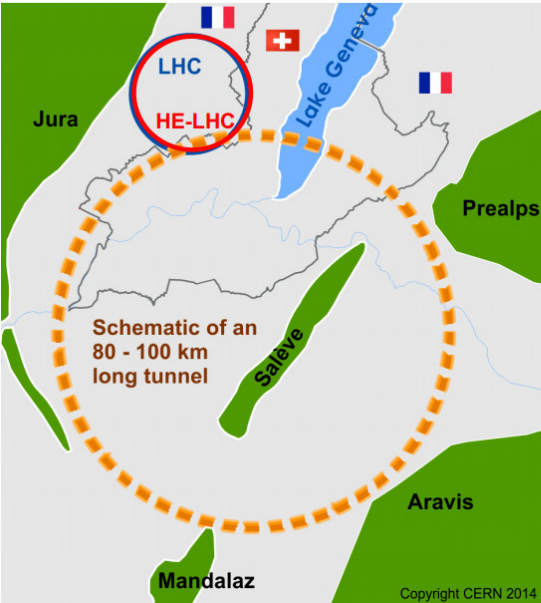
⇒ arXiv::1908.02302

What can happen in the future?



What can happen in the future?





International FCC collaboration with CERN as host lab to study:

- ~100 km tunnel infrastructure in Geneva area, linked to CERN
- e^+e^- collider (*FCC-ee*),
→ potential first step
- pp -collider (*FCC-hh*)
→ long-term goal, defining infrastructure requirements

~16 T \Rightarrow 100 TeV pp in 100 km

- **HE-LHC** with *FCC-hh* technology
- **Ions** and **lepton-hadron** options with hadron colliders

FCC-ee:

- **~20-50 fold improved precision on many EW quantities (equiv. to factor 5-7 in mass)** (m_Z , m_W , m_{top} , $\sin^2 \theta_W^{eff}$, R_b , $\alpha_{QED}(m_Z)$, $\alpha_s(m_Z)$, $\alpha_s(m_Z, m_W, m_\tau)$, Higgs and top quark couplings)
- **Exploring 10 - 100 TeV energy scale via couplings with precision measurements**
- **Machine design for highest luminosities at Z, WW, ZH and ttbar working points**

FCC-hh:

- **Highest center of mass energy for direct production up to 20 - 30 TeV**
- **Huge rates for single and multiple production of SM bosons (H,W,Z) and quarks**
- **Machine design for ~100 TeV c.m. energy & int. luminosity ~ 20ab⁻¹ in 25 years**

HE-LHC:

- **Doubling LHC collision energy with FCC-hh 16 T magnet technology**
- **c.m. energy ~ 27 TeV = 14 TeV x 16 T/8.33T, target luminosity ≥ 4 x HL-LHC**
- **Machine design within constraints from LHC CE and using HL-LHC and FCC techn.**

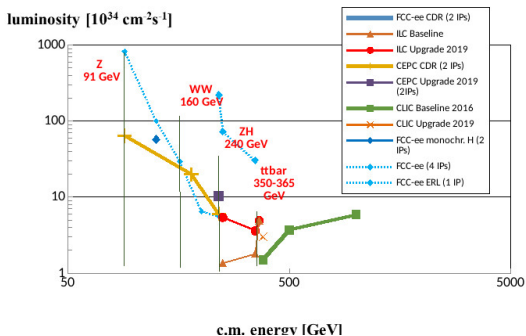
working point	assumed typical luminosity/IP [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$] = design value minus 15(10)%	total luminosity (2 IPs)/ yr; half of typical luminosity assumed in 1st two years (Z) and 1st year ($t\bar{t}$)	physics goal	run time [yr]
Z first 2 years	100	26 $\text{ab}^{-1}/\text{year}$	150 ab^{-1}	4
Z later	200	48 $\text{ab}^{-1}/\text{year}$		
W	25	6 $\text{ab}^{-1}/\text{year}$	10 ab^{-1}	1-2
H	7.0	1.7 $\text{ab}^{-1}/\text{year}$	5 ab^{-1}	3

machine modification for RF installation & rearrangement: 1 year

top 1st year (350 GeV)	0.8	0.2 $\text{ab}^{-1}/\text{year}$	0.2 ab^{-1}	1
top later (365 GeV)	1.4	0.34 $\text{ab}^{-1}/\text{year}$	1.5 ab^{-1}	4

total program duration: 15 years – incl. machine modifications
phase 1 (Z, W, H): 9 years, phase 2 (top): 6 years

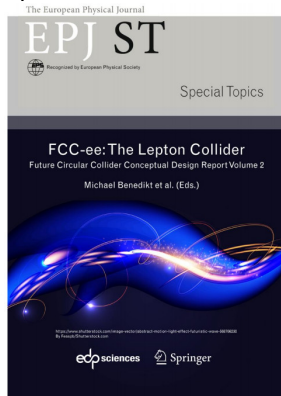
FCCee in context



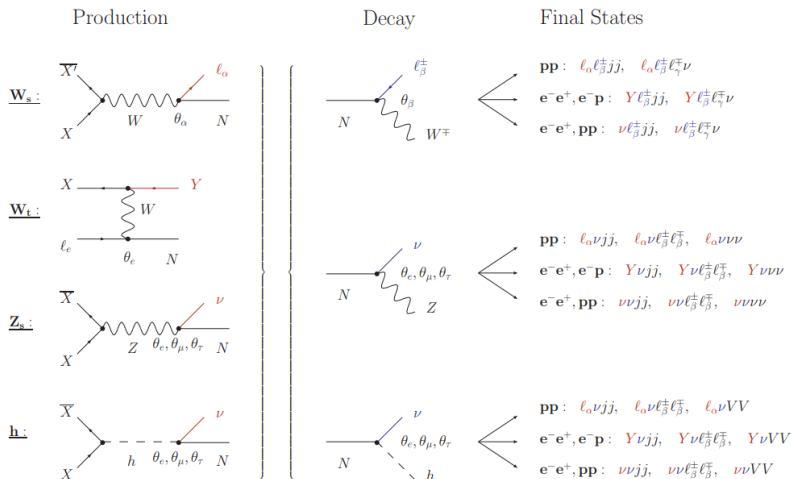
⇒ Check out the CDR : CERN-ACC-2018-0057

⇒ Also the theory report: arXiv:1905.05078

⇒ The FCCee is the most efficient machine up to the $t\bar{t}$ threshold.

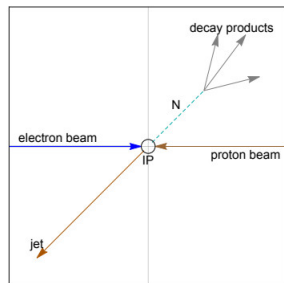
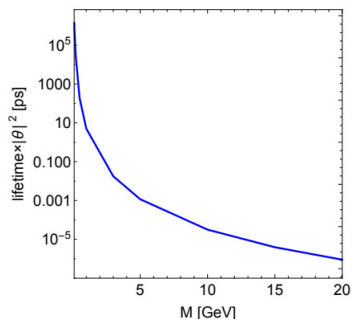


Schematizing sterile neutrino searches at FCC



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

Displaced vertexes

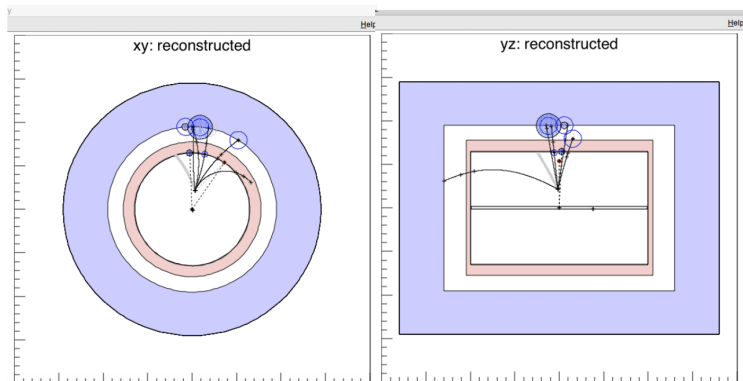


Example: FCC-he

- ⇒ In the interesting region: $m < m_W$ and $\theta < 10^{-5}$
- ⇒ Displacement: measurement of primary (production) vertex.
- ⇒ Secondary vertex with „large” displacement
- ⇒ ee he: A few times tracking resolution: $\mathcal{O}(10)\mu\text{m}$,
- ⇒ hh: Beyond background, detector noise, pileup: $\mathcal{O}(10)\text{cm}$.

Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

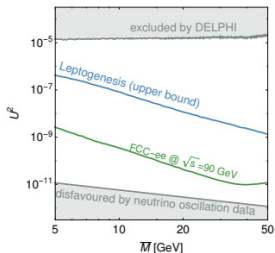
Displaced vertexes



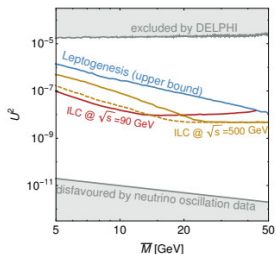
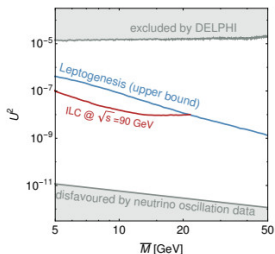
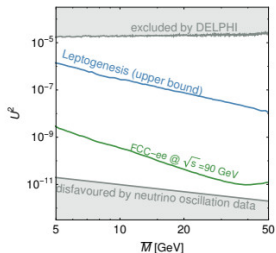
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Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.0272

Normal Ordering

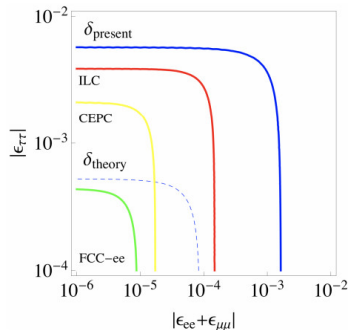


Inverted Ordering

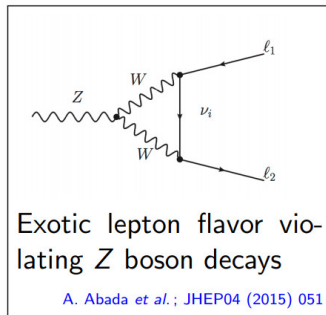


arXiv:1710.03744, S.Antusch, E.Cazzato, M.Drewes, O.Fischer, B.Garbrecht, D.Gueter, J.Klaric

FCCee indirect

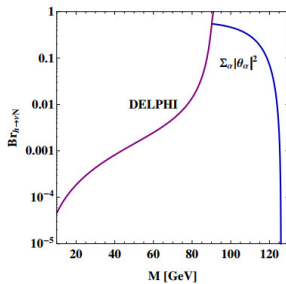
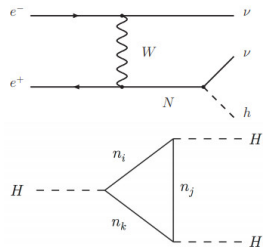


Antusch, OF; JHEP 1410 (2014) 094



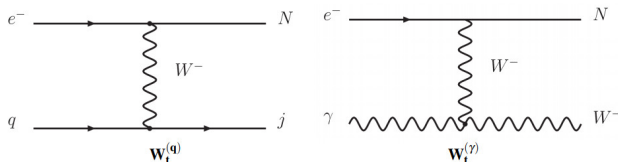
- ⇒ Modification of the theory prediction of precision observables.
- ⇒ Also CKM unitarity, cLFV, LUV.
- ⇒ Currently still dominated by LEP!

FCCee Higgs portal



S. Antusch, OF; JHEP 1604 (2016) 189

- ⇒ Mono-Higgs production mechanism!
- ⇒ New Higgs decays:
 - Modification of Higgs Branching fractions.
 - New decays: $H \rightarrow N\nu$.
 - Invisible width modification.
- ⇒ Modification of triple Higgs coupling.

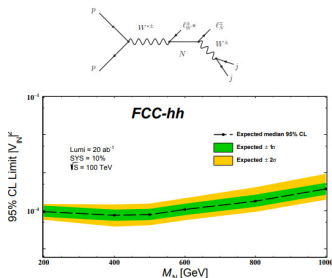


- ⇒ Large Lorentz boost makes the displaced vertexes clearly visible.
- ⇒ Many final states to look at:

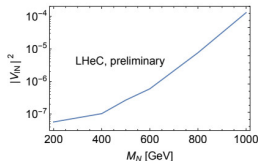
Name	Final State	$ \theta_\alpha $ Dependency	LFV
lepton-trijet	$jjj\ell_\alpha^-$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	✓
jet-dilepton	$j\ell_\alpha^-\ell_\beta^+\nu$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}^{(*)}$	✓
trijet	$jjj\nu$	$ \theta_e ^2$	×
monojet	$j\nu\nu\nu$	$ \theta_e ^2$	×

Name	Final State	$ \theta_\alpha $ Dependency	LFV
lepton-quadrjet	$jjjj\ell_\alpha^-$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}$	✓
dilepton-dijet	$\ell_\alpha^-\ell_\beta^+\nu jj$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}^{(*)}$	✓
trilepton	$\ell_\alpha^-\ell_\beta^-\ell_\gamma^+\nu\nu\nu$	$\frac{ \theta_e\theta_\alpha ^2}{\theta^2}^{(*)}$	✓
quadrjet	$jjjj\nu$	$ \theta_e ^2$	×
electron-di-b-jet	$e^-bb\nu\nu$	$ \theta_e ^2$	×
dijet	$jj\nu\nu\nu$	$ \theta_e ^2$	×
monolepton	$\ell_\alpha^-\nu\nu\nu\nu$	$ \theta_e ^2$	×

Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

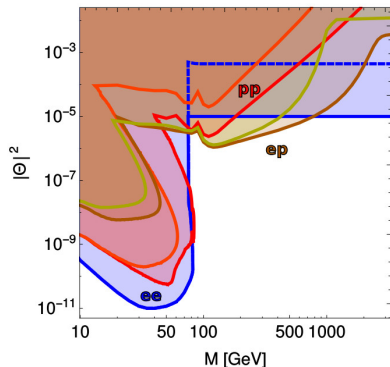


S. Antusch, E. Cazzato, O. Fischer, A. Hammad and K. Wang, JHEP 1810 (2018) 067



- ⇒ LFV is the thing to look for!!!
- ⇒ The best final states: $l_{\alpha}^{\pm} l_{\beta}^{\mp} j j$, $l^{\pm} l^{\mp} l_{\gamma}^{\pm}$
- ⇒ For ep machine the more sensitive ones are: $\mu j j j$ and $\tau j j j$.
- ⇒ Also LNU are there: $\mu^{\pm} \mu^{\pm} j$ (pp) and $e^{+} j$ (ep).

FCC in total



Credit to S.Antusch, E.Cazzato, O.Fischer, arXiv::1612.02728

⇒ FCCee:

- Dominates the exclusion below the m_W mass.
- Precision indirect constraints: EWPO, CKM, etc.

⇒ FCCeh, FCChh:

- Sensitivity in high mass region.
- Higgs potential.
- LFV, LNV.

Summary

- ⇒ Hunting for RHN is very well motivated.
- ⇒ Neutrino program has to be considered a core of future colliders.
- ⇒ FCC has unique sensitivity for RHN!
- ⇒ Huge amount of measurements and constraints to be performed.
- ⇒ Complementarity between different colliders.

Credit to M. Drewers, [Slides]

